

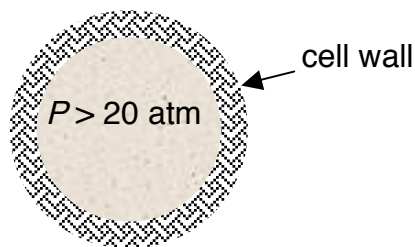
## PHYS 4xx Net 6 - Network percolation and failure

### Strain at failure

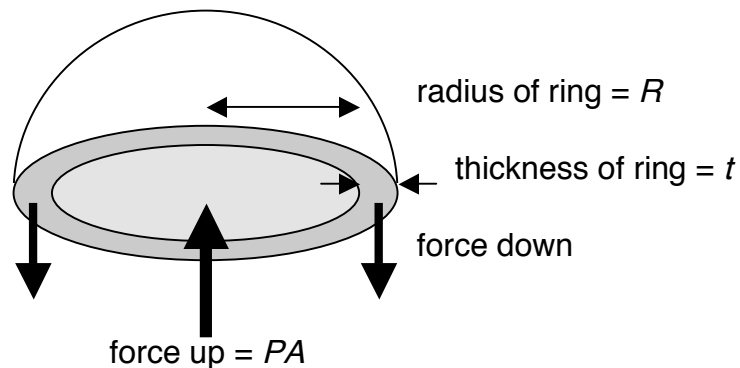
- strain at which a material fails varies widely
- many lipid bilayers fail at 2-5% strain
- some protein networks can withstand strains of 30% or more
- brittle materials fail at very small strains

### Stress at failure

Let's use a pressurized bacterium to provide a benchmark for the stresses some components of the cell must withstand.



Take a section through the cell, and applying a force balance equation at the cell wall:



Now, the force acting up across the area of the section is

$$F_{UP} = [\text{pressure}] \cdot [\text{area}]$$

or

$$F_{UP} = \pi R^2 P$$

where  $R$  is the radius of the cell. The force is **not**  $2\pi R^2 P$ .

The force acting down from the cell wall (which has the shape of a ring in cross section) has a similar form, with the 3D stress replacing the pressure in the fluid.

$$F_{DOWN} = [\text{3D stress}] \cdot [\text{area of the ring}]$$

If the thickness of the shell is small,  $t \ll R$ , then the area of the ring is  
 $[area] = [perimeter] \cdot [thickness] = 2\pi Rt$ .

Equating the forces gives

$$F_{DOWN} = F_{UP}$$

$$[3D \text{ stress}] \cdot 2\pi Rt = P \pi R^2$$

or

$$[3D \text{ stress}] t = PR / 2.$$

This is known as the law of Laplace, which was originally derived for soap bubbles.

OK, so let's put in some numbers.

radius of bacterium = 1  $\mu\text{m}$

[3D stress at failure] =  $2 \times 10^7 \text{ J/m}^3$  (typical value for many materials)

maximum pressure required = 20 atm =  $2 \times 10^6 \text{ J/m}^3$ .

Then the predicted thickness of the cell wall to carry this stress is

$$\begin{aligned} t &= PR / (2[\text{stress}]) \\ &= 2 \times 10^6 \times 10^{-6} / (2 \times 2 \times 10^7) \\ &= 50 \text{ nm} \end{aligned}$$

This is pretty close for Gram positive bacteria that have to withstand these pressures. Antibiotics take advantage of these high pressures to destroy bacteria.

### *Griffith formula*

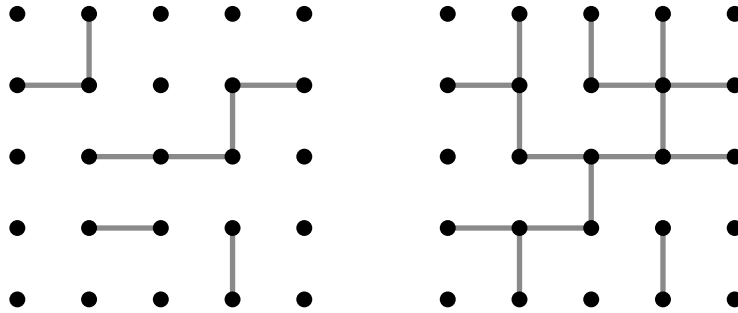
The failure of brittle materials has been well studied; one of the fundamental results is that the critical stress  $\sigma_c$  for the propagation of a crack is given by

$$\sigma_c \sim (2Y\gamma_s / \pi a)^{1/2},$$

where  $Y$  is the Young's modulus,  $\gamma_s$  is the interfacial tension between the exposed surface of the crack and the medium. The crack length is  $2a$ .

### **Percolation**

What happens if the chain density is so low that a connected lattice doesn't exist? Consider the two-dimensional square lattice, on which bonds have been placed randomly



$p$  = probability that a given bond site is occupied ( $p = 1$  for no vacancies)  
 $p$  is low on the left hand configuration, and no connected path crosses the lattice  
 $p$  is higher on the right hand side (but still  $< 1$ ) and a connected path exists  
 for infinite systems, a connecting path appears at the connectivity percolation threshold  $p_c$

$p_c = 0.5$  for a square lattice in 2D and  $p_c \sim 0.35$  for a triangular lattice in 2D

Feng and Sen (1984) showed that the elastic moduli vanish below a distinct rigidity percolation threshold  $p_R$ , which may be larger than  $p_c$  at zero temperature

One argument based on the counting of constraints gives (Maxwell, 1864)

number of bonds linked to a given vertex is, on average

$$[\text{no. of bonds connected per site}] = zp$$

where  $z$  is the number of links available

for  $N$  sites, the total number of bonds is

$$[\text{total no. bonds}] = Nzp/2$$

where the factor of 2 comes from each bond being shared by two sites

Now, there are  $Nd$  degrees of freedom for the lattice in a  $d$ -dimensional space, so the net number of degrees of freedom after applying the bond constraints is

$$[\text{no. of floppy modes}] = Nd - Nzp/2.$$

at the failure point, the number of floppy modes vanishes, so the occupation probability at failure  $p^*$  is

$$p^* = 2d/z$$

Measured value of  $\mu$  for red blood cells decreases with spectrin density; appears to be closer to  $p_c$  than  $p_R$

*Percolation threshold for a selection of regular lattices in two and three dimensions*

Lattice	$z$	$p_c$ (bond)	$p_c$ (site)	$p^*$ (bond)
<i>Two dimensions</i>				
honeycomb	3	0.653	0.696	1
square	4	0.500	0.593	1
triangular	6	0.347	0.500	2/3
<i>Three dimensions</i>				
simple cubic	6	0.249	0.312	1
body-centered cubic	8	0.180	0.246	3/4
face-centered cubic	12	0.119	0.198	1/2

"Bond" and "site" refer to different measures of defects in the lattice, as described in the text; Data from Chapter 2 of Stauffer and Aharony (1992).