

compact
layer

 $c_{\text{ox,red}}^{\text{s}}$
 $c_{\text{ox,red}}^{\text{b}}$

Note, that we explicitly distinguish concentrations at surface $c_{\text{red}}^{\text{s}}, c_{\text{ox}}^{\text{s}}$ from those in bulk $c_{\text{red}}^{\text{b}}, c_{\text{ox}}^{\text{b}}$. Concentrations in bulk (far away from the interface) determine E^{eq} .

Under equilibrium conditions:

$$c_{\text{red}}^{\text{s}} = c_{\text{red}}^{\text{b}}, \quad c_{\text{ox}}^{\text{s}} = c_{\text{red}}^{\text{b}}$$

No fluxes of species exist. We will see in the context of mass transport limitations why the distinction between surface and bulk concentrations is important!

At **finite current**, surface concentrations will generally be different from bulk concentrations, due to mass transport from bulk to surface:

$$c_{\text{red}}^{\text{s}} \neq c_{\text{red}}^{\text{b}}, \quad c_{\text{ox}}^{\text{s}} \neq c_{\text{red}}^{\text{b}}$$

Particular case:

mass transport from bulk to surface is fast

rate of diffusion >> rate of reaction

(see section on mass transport)

$$\Rightarrow c_{\text{red}}^{\text{s}} = c_{\text{red}}^{\text{b}} \text{ and } c_{\text{ox}}^{\text{s}} = c_{\text{ox}}^{\text{b}}$$

(note: not in equilibrium!)

\Rightarrow simplest variant of the **Butler-Volmer equation**

(historically the first version),

$$\text{(BV 3)} \quad j = j^0 \left\{ \exp\left(\frac{\alpha F \eta}{RT}\right) - \exp\left(-\frac{(1-\alpha) F \eta}{RT}\right) \right\}$$

Now: Three equations, each of them referred to as
“Butler-Volmer equation”

Which one should you use? – it depends!

General situation: complex, distributed electrode

(BV 3) is the least general form. It can only be used, when **all sorts of mass transport limitations can be neglected** (uniform distributions of reactants, simple electrodes, well-stirred solution) and current-potential relation is only determined by the kinetics of charge transfer.

Whenever you see people using (BV 3) in a situation with significant mass transport limitations, something is screwed.

(BV 2) can be used for **simple electrode structures**, when the **bulk concentrations are uniform** (i.e. E^{eq} is well-defined).

(BV 1) is the **most general form**. It explicitly displays all dependencies on concentrations and electrode potential. All other quantities in it are constants. This form is appropriate for complex distributed electrodes.

Pre-exponential factor in (BV 2) and (BV 3)

$$j^0 = Fk^0 c_{\text{red}}^{b(1-\alpha)} c_{\text{ox}}^{b\alpha}$$

is called exchange current density. It is an **important parameter** that characterizes a particular electrode reaction (Faradaic process). It is (like k^0) a measure of current density in anodic and cathodic directions **at zero overpotential** ($\eta = 0$), when the system is at equilibrium. It can be seen that it depends on the composition of the system (concentrations).

Physical meaning of j^0 : how does system respond to deviation of E from E^{eq} , i.e. applied $\eta \neq 0$?

At given E , which information about a system do you need to determine, in which direction the reaction will proceed and which current density it will generate?

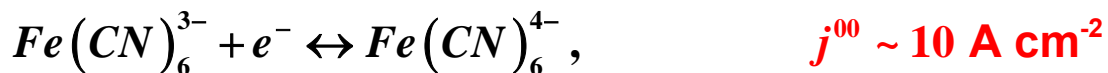
Definition of standard exchange current density:

$$j^{00} = Fk^0 \quad (\text{at standard concentration})$$

$$(\text{hence } j^0 = j^{00} (c_{\text{red}}^s)^{(1-\alpha)} (c_{\text{ox}}^s)^\alpha)$$

- values of j^{00} are reported in literature
(unless concentrations are explicitly specified)
- j^{00} is important property in electrode characterization

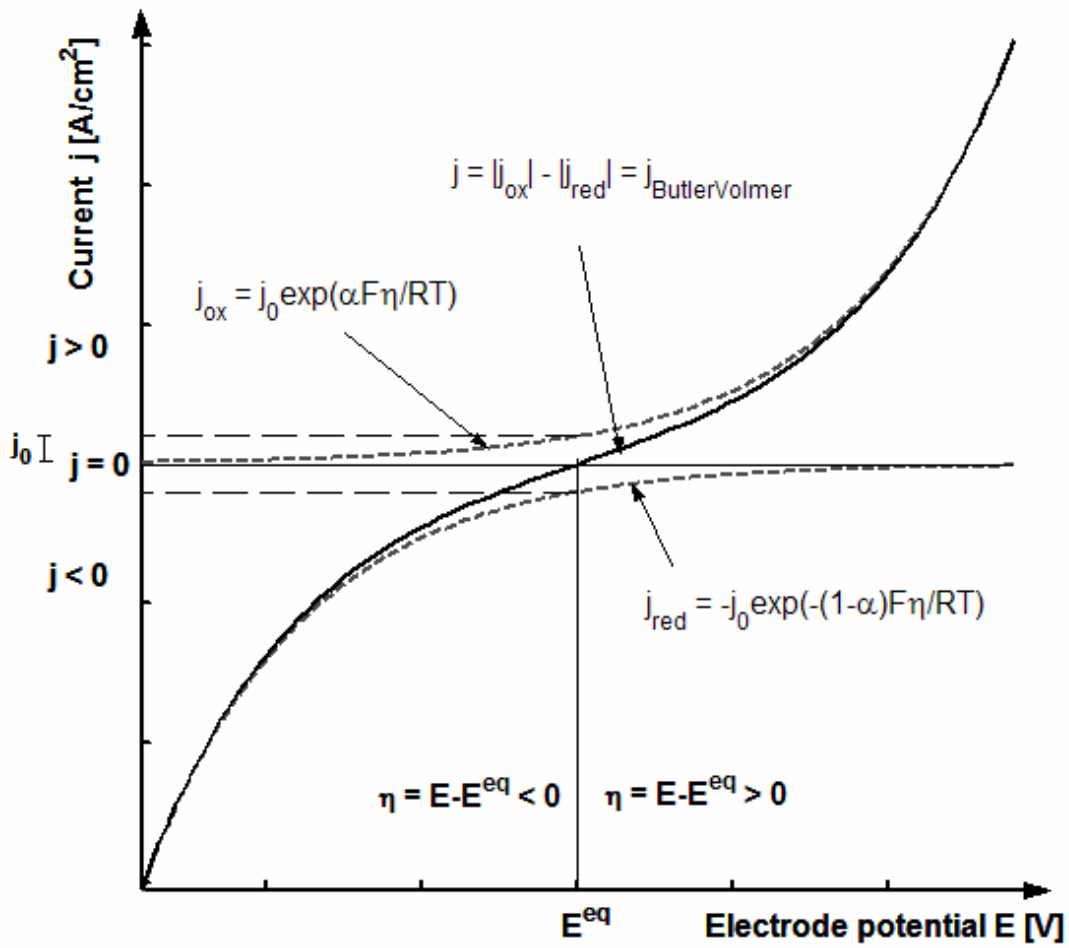
Values of j^{00} can widely differ for different reactions (Why?). In most cases, they are known only with low accuracy (within an order of magnitude). A few values:



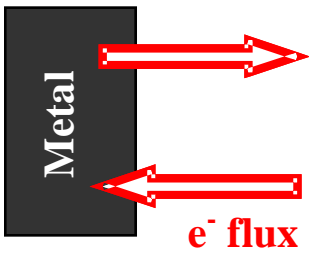
$j \ll j^0$: reversible electrode, Nernst-equation holds

ideally non-polarizable electrode: the potential will not change from equilibrium potential with the application of even a large current density

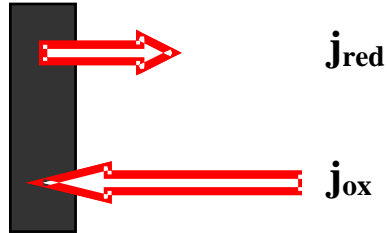
Meaning of the different parameters in BV equation:



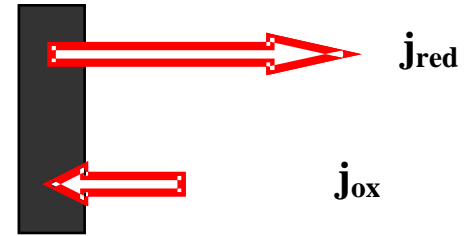
Three potential regions



$$E = E^{eq} ; \eta = 0, j = 0$$



$$E > E^{eq} ; \eta > 0 ; j > 0$$



$$E < E^{eq} ; \eta < 0 ; j < 0$$

1) **Exchange current density:** current density, which flows in each direction, anodic and cathodic, when the overall current is zero, i.e. in the electrochemical equilibrium. Hence, at E^{eq} we have a **dynamic equilibrium** where both oxidation and reduction take place at the **same rate**, given by j^0 .

As a pre-exponential factor, j^0 largely determines the overall kinetics:

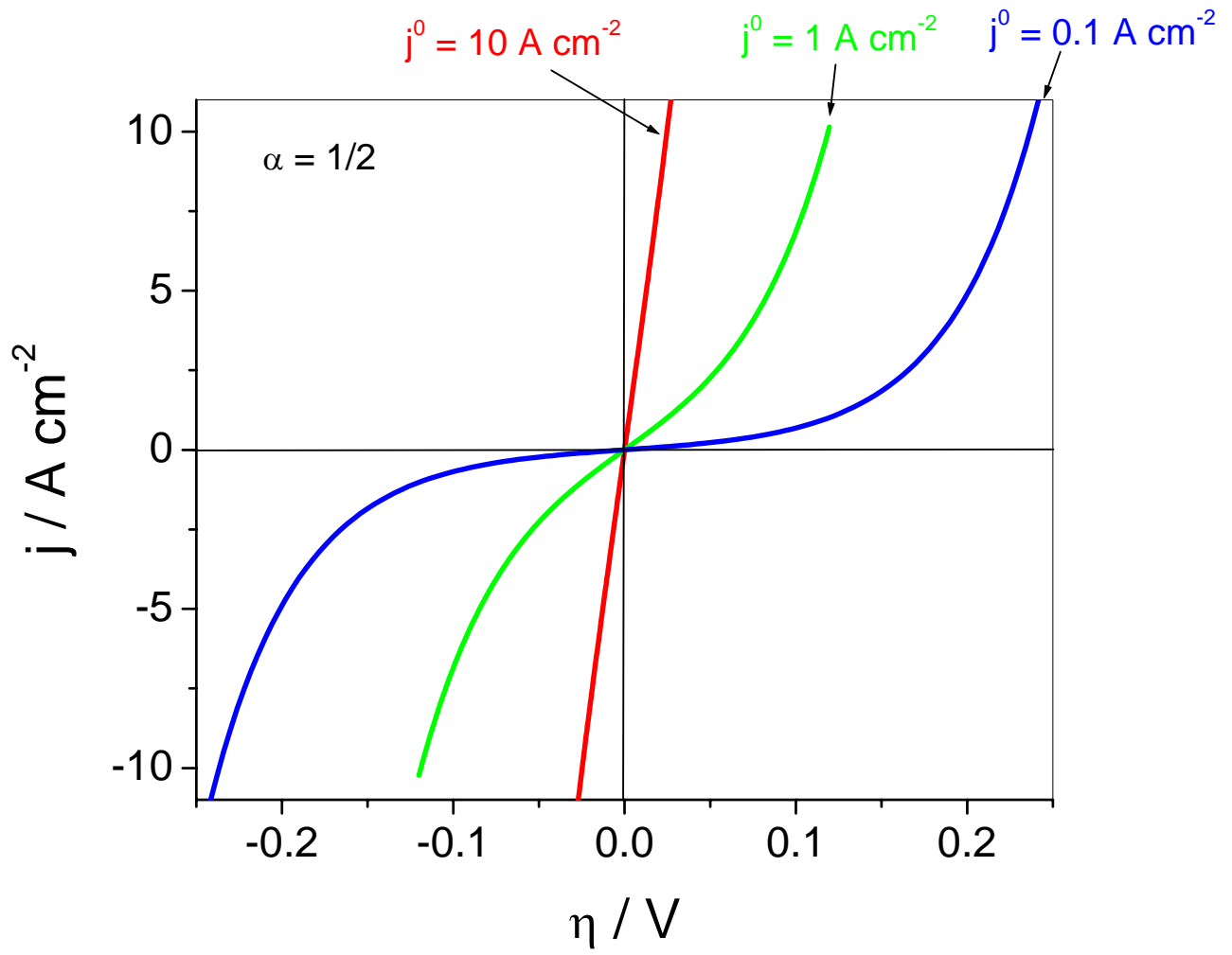
a) j^0 large (**fast electrode kinetics**):

High current densities reached at small $|\eta|$

b) j^0 small (**sluggish electrode kinetics**):

Large $|\eta|$ needed to obtain significant current density.

Effect of exchange current density

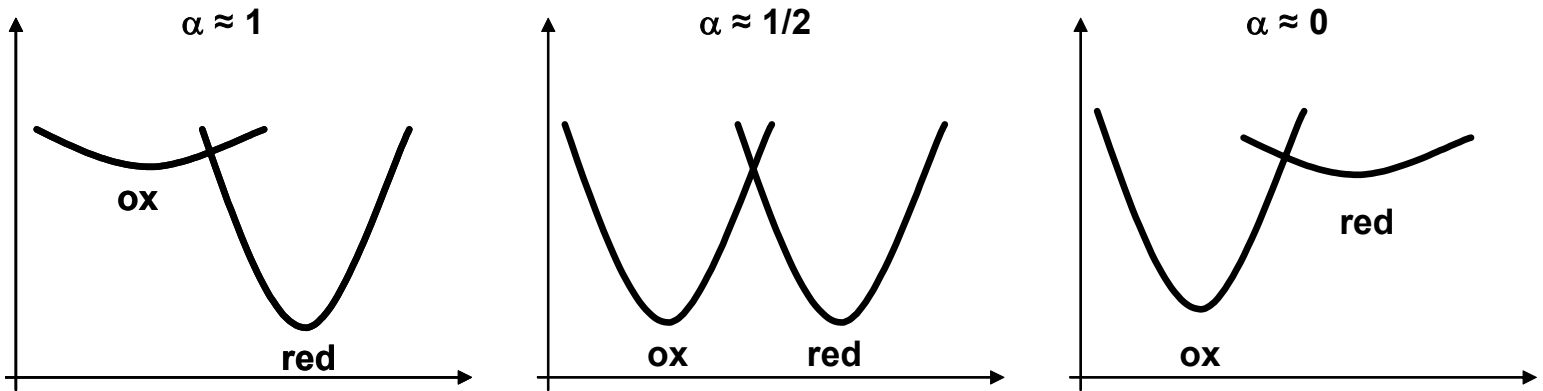


2) **Transfer coefficient α** : determines how strongly current density depends on the dependence of the on $|\eta|$ (often termed **symmetry factor**). Consider change $\Delta E = \Delta \eta$

- a) **If $\alpha = 0.5$** : j_{ox} increases by the same factor by which j_{red} decreases. The current density-overpotential graph is point-symmetric about $(E = E^{eq}, j = 0)$.
- b) **If $\alpha > 0.5$** : $|j|$ increases faster for oxidation than for reduction with the same $|\eta|$.
- c) **If $\alpha < 0.5$** : $|j|$ increases slower for oxidation than for reduction with the same $|\eta|$.

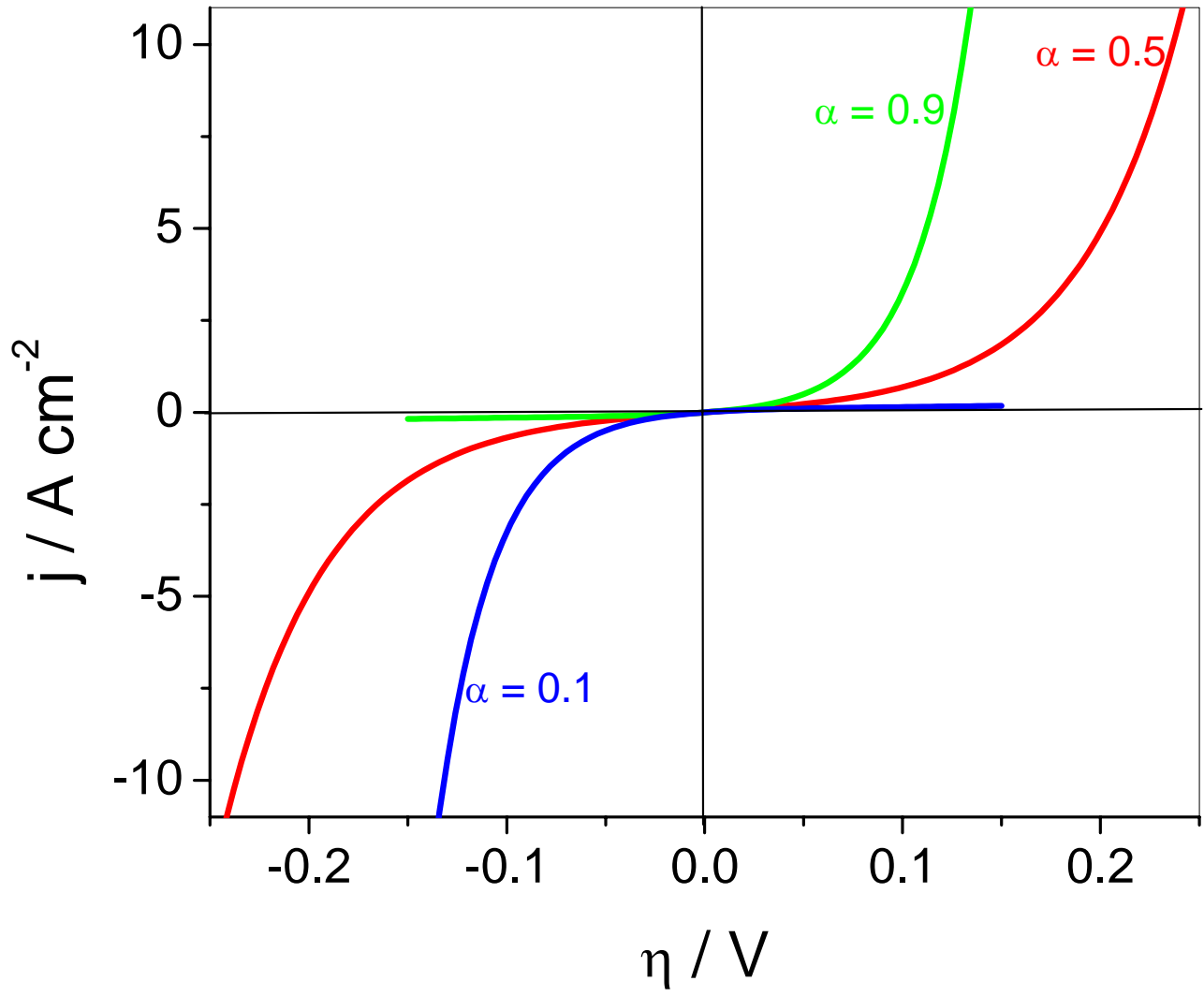
The following picture shows free energy profiles for different values of α

(anodic process: “ox” corresponds e^- in metal)



- usually (for metals) $\alpha \approx 1/2$
 - ⇒ symmetric e^- -transfer!
- activated complex has predominantly structure of oxidized state: $\alpha \approx 1$
 - ⇒ anodic reaction strongly dependent on E
- activated complex has predominantly structure of reduced state: $\alpha \approx 0$
 - ⇒ anodic reaction weakly dependent on E

Effect of transfer coefficient



3) Small overpotentials $|\eta| < \frac{RT}{F}$:

- negligible mass transport limitations - use (BV 3)
- both exponential terms important
- Taylor expansion in (BV 3) gives

$$j \approx j^0 \frac{F\eta}{RT} = \frac{\eta}{R_{ct}} \quad \text{where} \quad R_{ct} = \frac{RT}{j^0 F}$$

- linear relation between j and η
- form of Ohm's law with charge transfer resistance R_{ct} .

We can use the low overpotential regime to determine the **exchange current density!**

4) For large values of $|\eta|$, $|\eta| > 3 \frac{RT}{F}$ (the so-called **Tafel-regime**), one of the partial currents dominates so that the other one can be neglected in (BV 3). We still assume, that mass transport limitations are negligible.

a) $|\eta| > 3 \frac{RT}{F}$ and $\eta > 0 \Rightarrow j \approx j_{\text{ox}}$ and $j_{\text{red}} \approx 0$:

$$\log_{10}(j) = \log_{10}(j^0) + \frac{\alpha F}{2.3RT} \eta$$

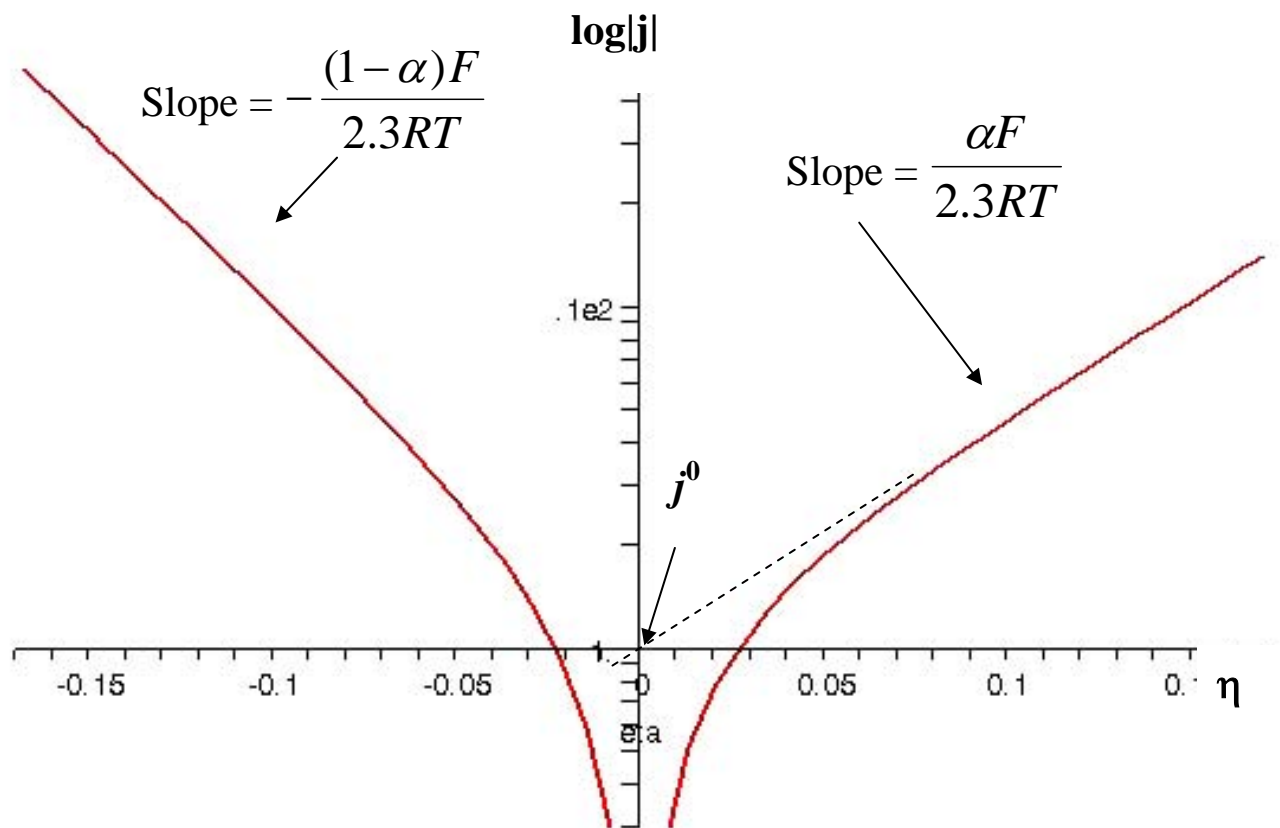
b) $|\eta| > 3 \frac{RT}{F}$ and $\eta < 0 \Rightarrow j \approx j_{\text{red}}$ and $j_{\text{ox}} \approx 0$:

$$\log_{10}(|j|) = \log_{10}(|j^0|) - \frac{(1-\alpha)F}{2.3RT} \eta$$

We can draw so-called **Tafel-plots: $\log|j|$ vs. η**

From such a plot, we can extract

- **transfer coefficient α** from the slope of a fit of the experimental data to a straight line at $|\eta| \gg 0$
- **exchange current density j^0** from the crossing point with the ordinate of this linear fit.



5) The inverse of the oxidation current slope in a Tafel plot is called the **Tafel-constant**:

$$b = \frac{2.3RT}{\alpha F}$$

Typical value?

6) At very large $|\eta|$, $|j|$ reaches a **limiting value** because the reaction rates is **limited by diffusion** of reactant to the reaction zone – see later lectures.

7) Most **charge transfer reactions** are in fact **inner-sphere reactions** and the **Butler-Volmer equation** does in general not apply. However, j still depends on η . It is convenient to retain the form of the Butler-Volmer equation with **empirical transfer coefficients** α and β :

$$K_{\text{ox}} = k^0 \exp\left\{\frac{\alpha F \eta}{RT}\right\} \text{ and } K_{\text{red}} = k^0 \exp\left\{-\frac{\beta F \eta}{RT}\right\}$$

In general: α and β depend on η and T and $\alpha + \beta \neq 1$