

Topics for Today

1.) The Black-Scholes Formula

- Assumptions + Notation
- Replicating Portfolio
- The No Arbitrage PDE
- Conversion to the 'Heat Equation'
- Using Fourier Transforms to Solve the Heat Equation
- Alternative Derivation I (Delta Hedging + "The Greeks")
- Alternative Derivation II (Risk-Neutral Pricing + Monte Carlo Simulation)
- Extensions + Complications
- Exotic Options
- Empirical Evaluation / Implied Volatility

Assumptions & Notation

- 1.) Underlying stock price governed by Geometric Brownian motion
- 2.) Riskless interest rate is constant
- 3.) Continuous-trading with no transactions costs
- 4.) No borrowing constraints or limits on short-selling
- 5.) Stock does not pay dividends (or produce other cash flows)
- 6.) No arbitrage possibilities!

Notation

S = Stock Price $(\frac{dS}{S} = \mu dt + \sigma dz)$

r = riskless rate $(\frac{dB}{B} = r dt)$

K = Strike Price

T = Expiration date

$C(S, t)$ = Price of call option

Replicating Portfolio

• Last time we derived a 1-period option pricing formula using a simple binomial tree argument. Now we extend to continuous-time.

• By Ito's Lemma we have,

$$dC = \frac{\partial C}{\partial S} dS + \frac{1}{2} \sigma^2 S^2 \frac{\partial^2 C}{\partial S^2} dt + \frac{\partial C}{\partial t} dt$$

$$= \left(\frac{\partial C}{\partial S} \mu S + \frac{1}{2} \frac{\partial^2 C}{\partial S^2} \sigma^2 S^2 + \frac{\partial C}{\partial t} \right) dt + \sigma S \frac{\partial C}{\partial S} \cdot dz$$

• Form a portfolio, P , consisting of x \$ in stock and y \$ in bonds, so that $P = x \cdot S + y \cdot B$. Therefore,

$$dP = x dS + y \cdot dB$$

$$= x [\mu S dt + \sigma S dz] + y r B dt$$

$$= (x \mu S + y r B) dt + x \sigma S \cdot dz$$

• Now pick x and y to replicate the option contract. Matching the dz terms we get

$$x = \frac{\partial C}{\partial S}$$

Next, pick y so that $P = \frac{\partial C}{\partial S} \cdot S + y B = C$. This gives,

$$y = \frac{1}{B} \left(C - \frac{\partial C}{\partial S} \cdot S \right)$$

The No Arbitrage PDE

- Plug $x + y$ into the previous expression for dP

$$\begin{aligned}dP &= (x\mu S + y r B) dt + x\sigma S \cdot dZ \\ &= \left[\frac{\partial C}{\partial S} \mu S + r \left(C - \frac{\partial C}{\partial S} \cdot S \right) \right] dt + \frac{\partial C}{\partial S} \sigma S \cdot dZ\end{aligned}$$

- Note that the dZ term is the same as the dZ term in dC . Matching the drift (dt) terms gives:

$$\frac{\partial C}{\partial S} \mu S + \frac{1}{2} \frac{\partial^2 C}{\partial S^2} \sigma^2 S^2 + \frac{\partial C}{\partial t} = \frac{\partial C}{\partial S} \mu S + r \left(C - \frac{\partial C}{\partial S} \cdot S \right)$$

- Re-arranging we get,

$$rC = \frac{\partial C}{\partial t} + rS \frac{\partial C}{\partial S} + \frac{1}{2} \sigma^2 S^2 \frac{\partial^2 C}{\partial S^2}$$

Black-Scholes
> PDE!

- This must be solved subject to the boundary conditions:

$$C(S, T) = \max\{S_T - K, 0\}$$

$$C(0, t) = 0$$

Conversion to the Heat Equation

• In general, PDEs do not have analytical solutions. However, it turns out that with appropriate transformations, the Black-Scholes PDE can be converted to the "Heat Equation", which is a famous PDE in mathematical physics with a well known solution.

• To see this, define the following change of variables:

$$x = \log\left(\frac{S}{K}\right)$$

$$\tau = \frac{1}{2}\sigma^2(T-t)$$

and define $C(S, t) = K \cdot V(x, \tau)$

• Note (using the chain rule)

$$\frac{\partial C}{\partial t} = -\frac{1}{2}\sigma^2 K \frac{\partial V}{\partial \tau}$$

$$\frac{\partial C}{\partial S} = \frac{K}{S} \frac{\partial V}{\partial x}$$

$$\frac{\partial^2 C}{\partial S^2} = \frac{K}{S^2} \left(\frac{\partial^2 V}{\partial x^2} - \frac{\partial V}{\partial x} \right)$$

• This implies the following equation for $V(x, \tau)$

$$\frac{\partial V}{\partial \tau} = \frac{\partial^2 V}{\partial x^2} + (k-1) \frac{\partial V}{\partial x} - k V$$

where $k = \frac{2r}{\sigma^2}$

- Note that the boundary condition at $t = T$ becomes an initial condition at $\tau = 0$

$$C(S, T) = \max(S - K, 0) = \max(Ke^x - K, 0)$$

Since $V = \frac{1}{K} C \Rightarrow \boxed{V(x, 0) = \max(e^x - 1, 0)}$ > New initial condition

- Next, let's change variables again, and define,

$$V = e^{\alpha x + \beta \tau} u(x, \tau)$$

$$\Rightarrow V_\tau = \beta e^{\alpha x + \beta \tau} u + e^{\alpha x + \beta \tau} u_\tau$$

$$V_x = \alpha e^{\alpha x + \beta \tau} u + e^{\alpha x + \beta \tau} u_x$$

$$V_{xx} = \alpha^2 e^{\alpha x + \beta \tau} u + 2\alpha e^{\alpha x + \beta \tau} u_x + e^{\alpha x + \beta \tau} u_{xx}$$

- This gives the following equation for u

$$u_\tau = u_{xx} + [2\alpha + (k-1)]u_x + [\alpha^2 + (k-1)\alpha - k - \beta]u$$

Note that if we choose $\alpha = -\frac{1}{2}(k-1)$ $\beta = -\frac{1}{4}(k+1)^2$

we get

$$\boxed{u_\tau = u_{xx}}$$

> Heat Equation

with transformed initial condition

$$u(x, 0) = \max \left[e^{\frac{1}{2}(k+1)x} - e^{\frac{1}{2}(k-1)x}, 0 \right]$$

Using Fourier Transforms to Solve the Heat Equation

- The basic idea behind all transform methods is the following:
 - 1.) Use the transform to convert a hard eq. into an easy eq.
 - 2.) Solve the easy equation
 - 3.) Invert the transform
- Given a function, $f(x)$, the Fourier transform and its inverse are:

$$\begin{aligned} \mathcal{F}[f] &= F(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\omega x} dx \\ \mathcal{F}^{-1}[F] &= f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(\omega) e^{i\omega x} d\omega \end{aligned} \quad \left. \vphantom{\begin{aligned} \mathcal{F}[f] &= F(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-i\omega x} dx \\ \mathcal{F}^{-1}[F] &= f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} F(\omega) e^{i\omega x} d\omega \end{aligned}} \right\} \begin{array}{l} \text{Fourier} \\ \text{Transform} \\ \text{pair} \end{array}$$

- Note the following

$$\mathcal{F}[f_x] = i\omega \mathcal{F}[f]$$

$$\mathcal{F}[f_{xx}] = -\omega^2 \mathcal{F}[f]$$

$$\mathcal{F}[f_c] = \frac{1}{c} \mathcal{F}[f]$$

$$\mathcal{F}[f * g] = \mathcal{F}[f] \mathcal{F}[g]$$

where $(f * g)(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x-s)g(s)ds$ → convolution

$$\Rightarrow \mathcal{F}^{-1}\{\mathcal{F}[f]\mathcal{F}[g]\} = f * g \quad \left. \vphantom{\Rightarrow \mathcal{F}^{-1}\{\mathcal{F}[f]\mathcal{F}[g]\} = f * g} \right\} \text{convolution Property}$$

• Letting $\hat{u} = \mathcal{F}[u]$ we then get

$$\frac{d\hat{u}}{d\tau} = -\omega^2 \hat{u} \quad \left. \vphantom{\frac{d\hat{u}}{d\tau}} \right\} \text{Fourier Transform of the Heat Equation}$$

Note, our PDE is now an ODE! The solution is:

$$\hat{u} = \hat{u}(0) e^{-\omega^2 \tau}$$

where $\hat{u}(0) = \mathcal{F}[u(x, 0)]$.

• Now invert using the Convolution Property,

$$u(x, \tau) = \mathcal{F}^{-1}[\hat{u}] = \mathcal{F}^{-1}[\hat{u}(0)] * \mathcal{F}^{-1}[e^{-\omega^2 \tau}]$$

• This implies,

$$u(x, \tau) = \frac{1}{\sqrt{4\pi\tau}} \int_{-\infty}^{\infty} u(s, 0) e^{-\frac{(x-s)^2}{4\tau}} ds$$

> Solution of the Heat Eq.

Comment

1.) The function $e^{-\frac{(x-s)^2}{4\tau}}$ goes by various names. Physicists call it a 'heat kernel'. Mathematicians call it a "Green's function". Economists refer to it as an "impulse response function".

The Black-Scholes Formula

- To finally get the Black-Scholes formula we just need to unwind all the transformations we did. This gives

$$C(S, t) = S \cdot \Phi \left[\frac{\log(S/K) + (r + \sigma^2/2)(T-t)}{\sigma \sqrt{T-t}} \right] - K e^{-r(T-t)} \Phi \left[\frac{\log(S/K) + (r - \sigma^2/2)(T-t)}{\sigma \sqrt{T-t}} \right]$$

where $\Phi(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^x e^{-s^2/2} ds$ is the Normal CDF.

- This is usually written as:

$$C(S, t) = S \cdot \Phi(d_1) - K e^{-r(T-t)} \cdot \Phi(d_2)$$

} Black
Scholes
Formula!

where $d_1 = \frac{\log(S/K) + (r + \sigma^2/2)(T-t)}{\sigma \sqrt{T-t}}$

$$d_2 = d_1 - \sigma \sqrt{T-t}$$

- Note $(d_1, d_2) \rightarrow +\infty$ as $t \rightarrow T$ if $S > K$ and $\rightarrow -\infty$ if $S < K$, which implies

$$C(S, T) = \max\{S - K, 0\} \quad \text{check } \checkmark$$

Alternative Derivation I

- Black + Scholes actually derive their PDE using a different (but equivalent) portfolio strategy.
- Rather than form a portfolio of the stock + bond which replicates the option contract, they form a portfolio consisting of the stock + option which is riskless.
- As before, let $C(S,t)$ be the value of the option, and let $P = \theta_1 C(S,t) + \theta_2 S$ be the value of a portfolio consisting of θ_1 unit of the option and θ_2 shares of stock.

• Again, from Ito's Lemma,

$$dP = \theta_1 [C_t + C_s ds + \frac{1}{2} \sigma^2 s^2 C_{ss}] dt + \theta_2 ds$$

- Suppose we set $\theta_1 = 1$ and $\theta_2 = -C_s$. Note that the terms involving ds cancel, and we get

$$dP = (C_t + \frac{1}{2} \sigma^2 s^2 C_{ss}) dt$$

- Note, this is non-random and so is riskless. Therefore to avoid arbitrage we must have

$$\begin{aligned} dP = rP dt &= r(C - S \cdot C_s) dt \\ &= C_t + \frac{1}{2} \sigma^2 s^2 C_{ss} \end{aligned}$$

$$\Rightarrow rC = C_t + rS C_s + \frac{1}{2} \sigma^2 s^2 C_{ss}$$

> This is the same PDE as before!

Delta Hedging and "The Greeks"

- The idea of going long one call + short $\frac{\partial C}{\partial S}$ shares of stock to form a riskless portfolio is called "delta hedging", where $\Delta = \frac{\partial C}{\partial S}$.
- Note, we could equivalently go long Δ shares of stock and short one call.
- More generally, the sensitivities of $C(S,t)$ to market conditions can be useful for hedging purposes.
- We have,

$$\begin{aligned} dC &= \frac{\partial C}{\partial S} ds + \frac{1}{2} \frac{\partial^2 C}{\partial S^2} (ds)^2 + \frac{\partial C}{\partial t} dt \\ &= \underbrace{\Delta}_{\text{Delta}} \cdot ds + \frac{1}{2} \underbrace{\Gamma}_{\text{Gamma}} \cdot (ds)^2 + \underbrace{\Theta}_{\text{Theta}} \cdot dt \end{aligned}$$

where (Δ, Γ, Θ) can be derived from the BS formula (e.g., $\Delta = N(d_1)$).

- Collectively, traders refer to these hedge ratios as "The Greeks". They tell you how the value of the option should change as a function of time and the underlying share price.