

Lecture 7. Wed. Sept. 22, 2004.

- Electric potential V — scalar!

$$V = \frac{U}{q} \quad U \text{ — potential energy.}$$

unit: Volt = J/C.

- Relation with \vec{E} :

① Direction of $\vec{E} \perp$ equipotential lines.

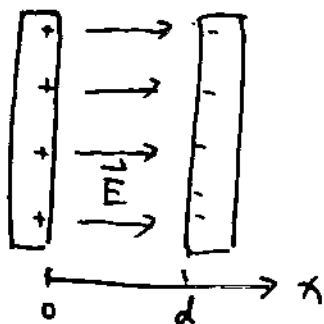
② $E = -\frac{\Delta V}{\Delta s}$ when Δs is along \vec{E} -line.

In general: $E_x = -\frac{\Delta V}{\Delta x}$

$$E_y = -\frac{\Delta V}{\Delta y}$$

$$E_z = -\frac{\Delta V}{\Delta z}$$

- e.g. Electric potential in a parallel-plate capacitor.



$$\vec{E} = E \hat{x} = E_x \hat{x}$$

from Gauss's Law: $E = \frac{\sigma}{\epsilon_0}$ (eq. 19-3)

σ — surface charge density.

$$\left(\sigma = \frac{q}{A} \right)$$

Then, along x -axis: $\Delta s = \Delta x$.

$$E = E_x = -\frac{\Delta V}{\Delta x} = \frac{\sigma}{\epsilon_0} \text{ (constant)}$$

$$\Delta V = -\frac{\sigma}{\epsilon_0} \cdot \Delta x$$

$$V_+ - V_- = -\frac{\sigma}{\epsilon_0} (0 - d)$$

Potential difference \equiv Voltage.

$$\boxed{V = \frac{\sigma d}{\epsilon_0}}$$

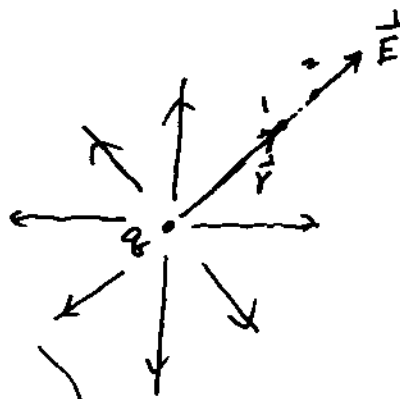
Voltage across a
parallel-plate capacitor.

e.g. electric potential of a point charge.

$\vec{E} \parallel \vec{r}$ always.

along a radius (along \vec{E} -line).

from point 1 to point 2.



$$E = -\frac{\Delta V}{\Delta r} \quad \left(\begin{array}{l} \text{only because} \\ \vec{r} \text{ is along } \vec{E}! \end{array} \right)$$

$$= -\frac{V_2 - V_1}{r_2 - r_1}$$

Then:

$$\Delta V = -E \cdot \Delta r, \quad dV = -E dr, \quad E = -\frac{dV}{dr}$$

$$V = -\int E dr$$

$$E = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2}$$

$$= -\int \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} dr$$

$$V = \frac{1}{4\pi\epsilon_0} \frac{q}{r} + C$$

Choose $V=0$ at ∞ (We often do this
for finite charge distributions)

Then $C=0$.

$$V = \frac{q}{4\pi\epsilon_0 r} = \frac{kq}{r}$$

electric potential
of a point charge
assuming $V(\infty)=0$.

• superposition principle of electric potential.

(Now it's algebraic sum!)

NOT vector sum.

V — scalar!