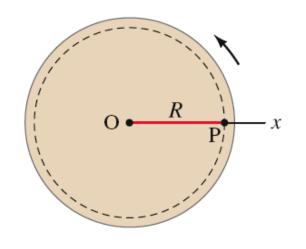
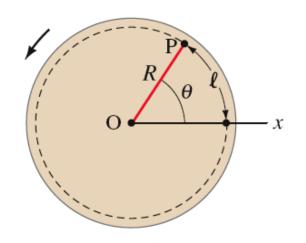
Phys101 Lectures 16, 17, 18 Rotational Motion

Key points:

- Rotational Kinematics
- Rotational Dynamics; Torque and Moment of Inertia
- Rotational Kinetic Energy
- Angular Momentum and Its Conservation

Ref: 8-1,2,3,4,5,6,8,9.





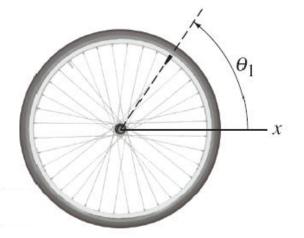
In purely rotational motion, all points on the object move in circles around the axis of rotation ("O"). The radius of the circle is R. All points on a straight line drawn through the axis move through the same angle in the same time. The angle θ in radians is defined:

$$\theta = \frac{l}{R},$$
 where l is the arc length.

Sign convention: + CCW

- CW

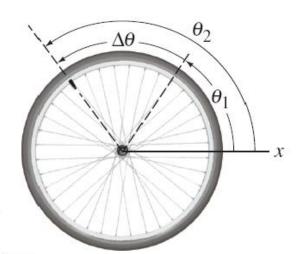
Angular displacement:



$$\Delta\theta = \theta_2 - \theta_1$$
.

The average angular velocity is defined as the total angular displacement divided by time:

$$\overline{\omega} = \frac{\Delta \theta}{\Delta t}.$$



The instantaneous angular velocity:

$$\omega = \lim_{\Delta t \to 0} \frac{\Delta \theta}{\Delta t} = \frac{d\theta}{dt}.$$

The angular acceleration is the rate at which the angular velocity changes with time:

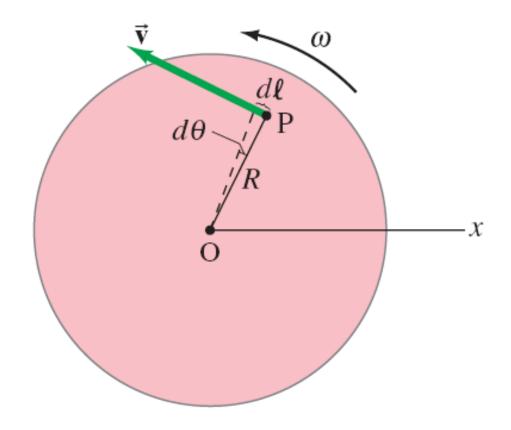
$$\overline{\alpha} = \frac{\omega_2 - \omega_1}{\Delta t} = \frac{\Delta \omega}{\Delta t}.$$

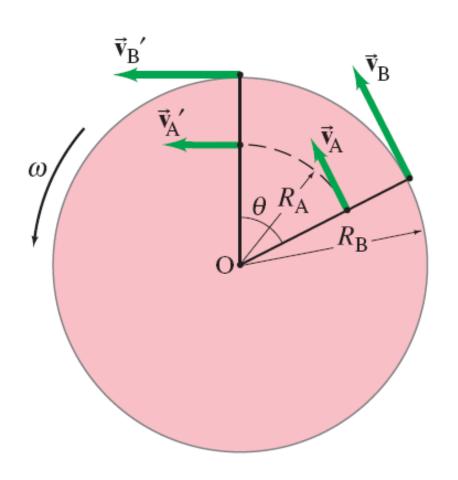
The instantaneous acceleration:

$$\alpha = \lim_{\Delta t \to 0} \frac{\Delta \omega}{\Delta t} = \frac{d\omega}{dt}$$

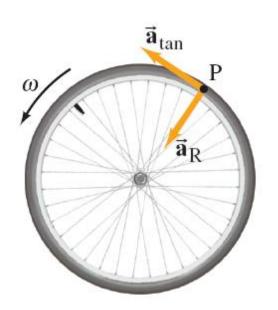
Every point on a rotating body has an angular velocity ω and a linear velocity ν .

They are related: $v = R\omega$.





Objects farther from the axis of rotation will move faster.



If the angular velocity of a rotating object changes, it has a tangential acceleration:

$$a_{\tan} = \frac{dv}{dt} = R \frac{d\omega}{dt} = R\alpha.$$

Even if the angular velocity is constant, each point on the object has a centripetal acceleration:

$$a_{\rm R} = \frac{v^2}{R} = \frac{(R\omega)^2}{R} = \omega^2 R.$$

Here is the correspondence between linear and rotational quantities:

TABLE 8-1 Linear and Rotational Quantities

Linear	Туре	Rotational	Relation
X	displacement	heta	$x = r\theta$
v	velocity	ω	$v = r\omega$
a_{tan}	acceleration	α	$a_{\tan} = r\alpha$

Copyright © 2005 Pearson Prentice Hall, Inc.

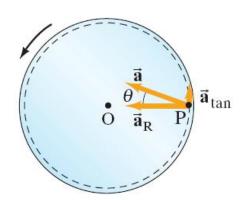
x represents arc length

$$a_{\rm R} = \frac{v^2}{R} = \frac{(R\omega)^2}{R} = \omega^2 R.$$



Example: Angular and linear velocities and accelerations.

A carousel is initially at rest. At t=0 it is given a constant angular acceleration $\alpha=0.060$ rad/s², which increases its angular velocity for 8.0 s. At t=8.0 s, determine the magnitude of the following quantities: (a) the angular velocity of the carousel; (b) the linear velocity of a child located 2.5 m from the center; (c) the tangential (linear) acceleration of that child; (d) the centripetal acceleration of the child; and (e) the total linear acceleration of the child.



The frequency is the number of complete revolutions per second:

$$f = \frac{\omega}{2\pi}$$
.

Frequencies are measured in hertz:

$$1 \text{ Hz} = 1 \text{ s}^{-1}$$
.

The period is the time one revolution takes:

$$T = \frac{1}{f}.$$

Constant Angular Acceleration

The equations of motion for constant angular acceleration are the same as those for linear motion, with the substitution of the angular quantities for the linear ones.

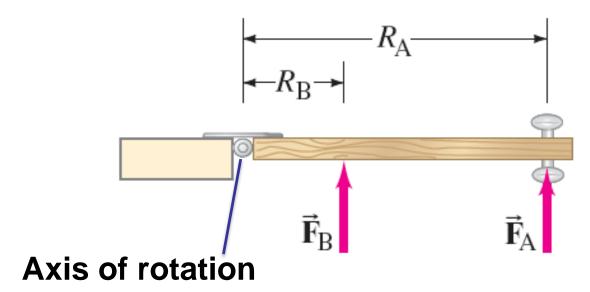
Angular	Linear		
$\omega = \omega_0 + \alpha t$	$v = v_0 + at$		
$\theta = \omega_0 t + \frac{1}{2}\alpha t^2$	$x = v_0 t + \frac{1}{2}at^2$		
$\omega^2 = \omega_0^2 + 2\alpha\theta$	$v^2 = v_0^2 + 2ax$		
$\overline{\omega} = \frac{\omega + \omega_0}{2}$	$\overline{v} = \frac{v + v_0}{2}$		

Example: Centrifuge acceleration.

A centrifuge rotor is accelerated from rest to 20,000 rpm in 30 s. (a) What is its average angular acceleration? (b) Through how many revolutions has the centrifuge rotor turned during its acceleration period, assuming constant angular acceleration?

To make an object start rotating, a force is needed; both the position and direction of the force matter.

The perpendicular distance from the axis of rotation to the line along which the force acts is called the lever arm.



In a FBD, don't shift the forces sideways!



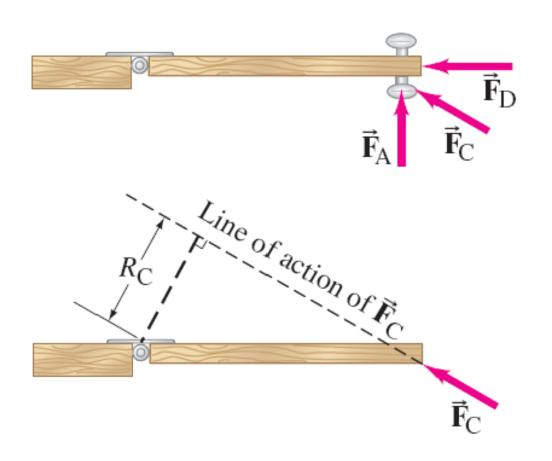
Axis of rotation

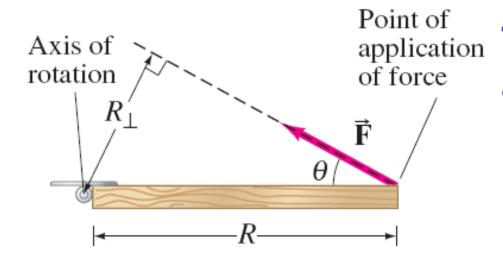


Axis of rotation

A longer lever arm is very helpful in rotating objects.

Here, the lever arm for $F_{\rm A}$ is the distance from the knob to the hinge; the lever arm for $F_{\rm D}$ is zero; and the lever arm for $F_{\rm C}$ is as shown.

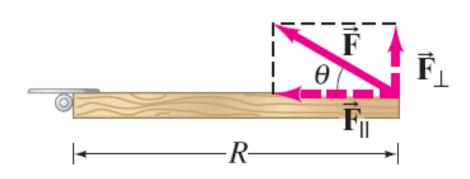




The torque is defined as:

$$\tau = R_{\perp} F$$
.

$$R_{\perp} = R \sin \theta$$



Also:

$$\tau = R F_{\perp}$$

That is,

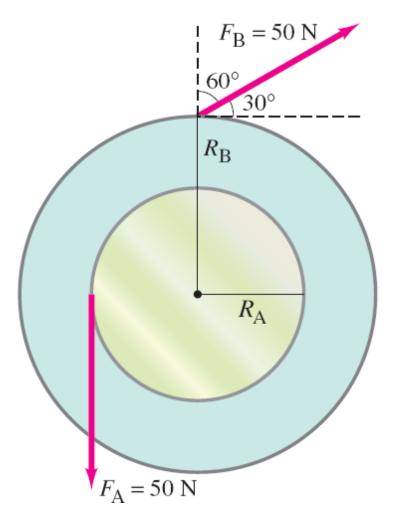
$$\tau = RF\sin\theta$$

$$F_{\perp} = F \sin \theta$$

 θ is the angle between \vec{F} and \vec{R} .

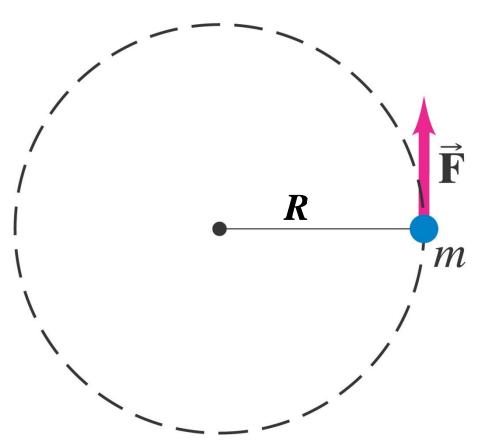
Example: Torque on a compound wheel.

Two thin disk-shaped wheels, of radii $R_{\rm A}$ = 30 cm and $R_{\rm B}$ = 50 cm, are attached to each other on an axle that passes through the center of each, as shown. Calculate the net torque on this compound wheel due to the two forces shown, each of magnitude 50 N.



Rotational Dynamics; Torque and Rotational Inertia

Knowing that F = ma, we see that $\tau = mR^2\alpha$.



$$\tau = FR = maR = mR^{2}\alpha$$
(since $a = \alpha R$)

This is for a single point mass; what about an extended object?

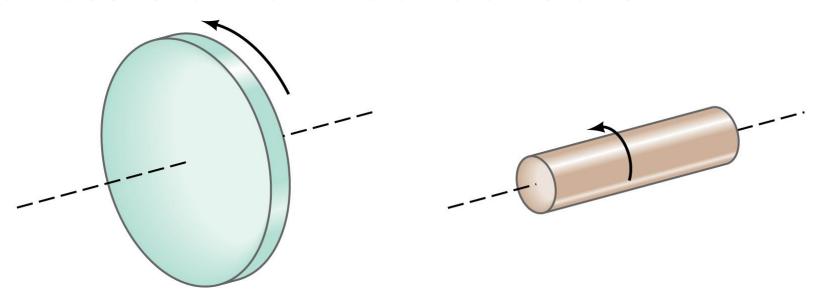
As the angular acceleration is the same for the whole object, we can write:

$$\Sigma \tau = (\Sigma mR^2)\alpha.$$

Rotational Dynamics; Torque and Rotational Inertia

The quantity $I = \sum m_i R_i^2$ is called the rotational inertia (moment of inertia) of an object.

The distribution of mass matters here—these two objects have the same mass, but the one on the left has a greater rotational inertia, as so much of its mass is far from the axis of rotation.

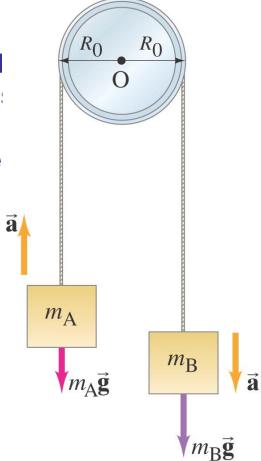


	Object	Location of axis		Moment of inertia
(a)	Thin hoop, radius R_0	Through center	Axis	MR_0^2
(b)	Thin hoop, radius R_0 width w	Through central diameter	Axis	$\frac{1}{2}MR_0^2 + \frac{1}{12}Mt$
(c)	Solid cylinder, radius R_0	Through center	Axis	$\frac{1}{2}MR_0^2$
(d)	Hollow cylinder, inner radius R_1 outer radius R_2	Through center	Axis R_2	$\frac{1}{2}M(R_1^2 + R_2^2)$
(e)	Uniform sphere, radius r_0	Through center	Axis	$\frac{2}{5}Mr_0^2$
(f)	Long uniform rod, length ℓ	Through center	Axis	$\frac{1}{12}M\ell^2$
(g)	Long uniform rod, length ℓ	Through end	Axis	$\frac{1}{3}M\ell^2$
(h)	Rectangular thin plate, length ℓ , width w	Through center	Axis	$\frac{1}{12}M(\ell^2+w^2)$

Rotational Dynamics; Torque and Rotational Inertia

The rotational inertia of an object depends not only on its mass distribution but also the location of the axis of rotation—compare (f) and (g), for example. **Example: Atwood's machine.**

An Atwood machine consists of two masses, $m_{\rm A}$ and $m_{\rm B}$, which are connected by a cord of negligible mass that passes over a pulley. If the pulley has radius R_0 and moment of inertia I about its axle, determine the acceleration of the masses $m_{\rm A}$ and $m_{\rm B}$.



Rotational Kinetic Energy

The kinetic energy of a rotating object is given by

$$K = \sum \left(\frac{1}{2}mv^2\right).$$

By substituting the rotational quantities, we find that the rotational kinetic energy can be written as:

rotational
$$K = \frac{1}{2}I\omega^2$$
.

$$K = \sum \left(\frac{1}{2}mv^2\right) = \frac{1}{2}\sum (mR^2\omega^2) = \frac{1}{2}\omega^2\sum (mR^2) = \frac{1}{2}\omega^2I$$

An object in both translational and rotational motion has both translational and rotational kinetic energy:

$$K = \frac{1}{2}Mv_{\rm CM}^2 + \frac{1}{2}I_{\rm CM}\omega^2$$
.

Example (Old final exam, #19):

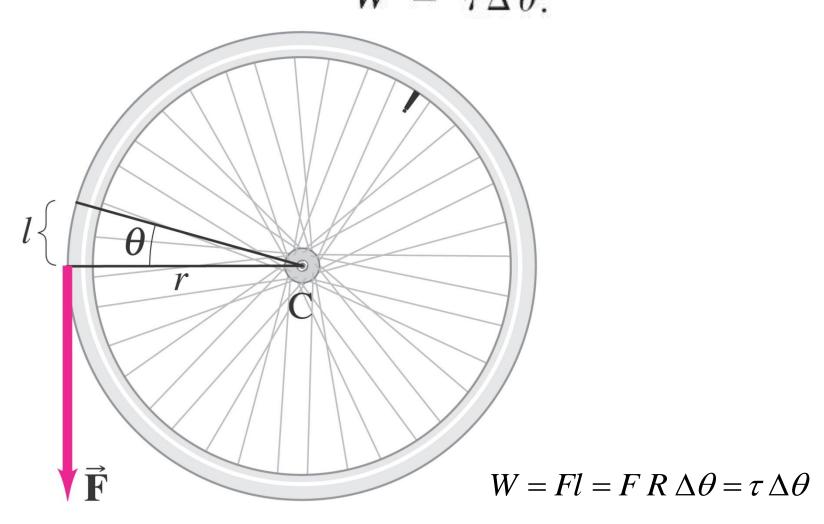
A pencil, 16 cm long (l=0.16m), is released from a vertical position with the eraser end resting on a table. The eraser does not slip. Treat the pencil like a uniform rod.

- (a) What is the angular acceleration of the pencil when it makes a 30° angle with the vertical?
- (b) What is the angular speed of the pencil when it makes a 30° angle with the vertical?

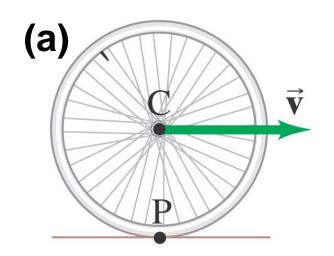
CM

Rotational Work

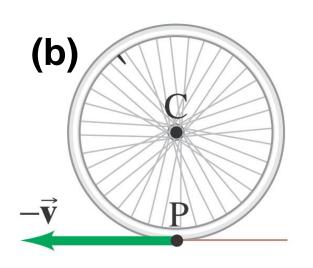
The torque does work as it moves the wheel through an angle θ : $W = \tau \Delta \theta$.



Rotational Plus Translational Motion; Rolling



In (a), a wheel is rolling without slipping. The point P, touching the ground, is instantaneously at rest, and the center moves with velocity \vec{v} .

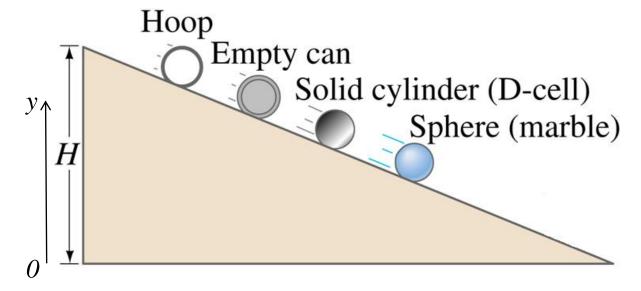


In (b) the same wheel is seen from a reference frame where C is at rest. Now point P is moving with velocity $-\vec{v}$.

The linear speed of the wheel is related to its angular speed:

$$v = R\omega$$
.

Demo: Which one reaches the bottom first? Why?



Angular Momentum—Objects Rotating About a Fixed Axis

The rotational analog of linear momentum is angular momentum, L:

$$L = I\omega$$
.

Then the rotational analog of Newton's second law is:

$$\Sigma \tau = \frac{dL}{dt}.$$

This form of Newton's second law is valid even if *I* is not constant.

Conservation of Angular Momentum

In the absence of an external torque, angular momentum is conserved:

$$\frac{dL}{dt} = 0$$
 and $L = I\omega = \text{constant}$.

The total angular momentum of a system remains constant if the net external torque is zero.

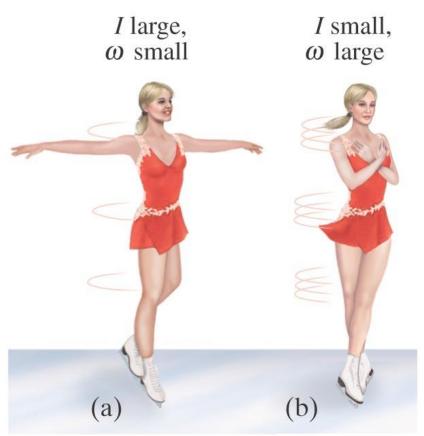
This means:

$$I_i \omega_i = I_f \omega_f = constant$$

Therefore, a figure skater can increase her angular speed by reducing her moment of inertia.

$$\omega_f = \frac{I_i}{I_f} \omega_i$$

Demo



i-clicker question 18-1

The condition for angular momentum to be conserved is

- (A) It's a closed system.
- (B) The net external force is zero.
- (C) No non-conservative work.
- (D) The net external torque is zero.
- (E) The angular momentum is always conserved.

Example: Object rotating on a string of changing length.

A small mass m attached to the end of a string revolves in a circle on a frictionless tabletop. The other end of the string passes through a hole in the table. Initially, the mass revolves with a speed $v_1 = 2.4$ m/s in a circle of radius $R_1 = 0.80$ m. The string is then pulled slowly through the hole so that the radius is reduced to $R_2 = 0.48$ m. What is the speed, v_2 , of the mass now?

Angular Momentum of a Moving Particle

You could open a door by throwing a large piece of playdoh at the door.

When $\vec{p} \perp \vec{R}$,

$$L = pR = mvR = m(R\omega)R = mR^2\omega = I\omega$$

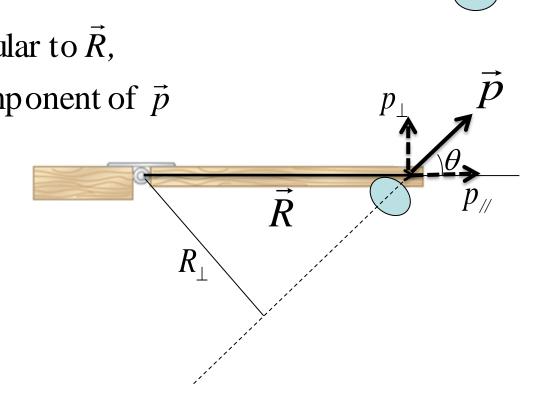
When \vec{p} is not perpendicular to \vec{R} , only the perpendicular component of \vec{p} contributes to L:

$$L = p_{\perp}R = pR \sin \theta$$

That is equivalent to

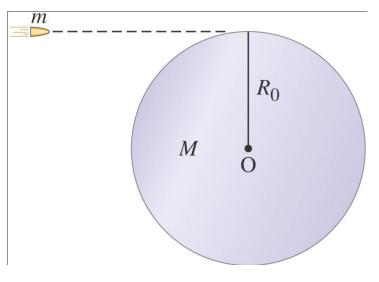
$$L = pR_{\perp} = pR \sin \theta$$

i.e.,
$$L = mvRsin\theta$$



Example: A bullet of mass m moving with velocity v strikes and becomes embedded at the edge of a cylinder of mass M and radius R_0 . The cylinder, initially at rest, begins to rotate about its symmetry axis, which remains fixed in position. Assuming no frictional torque, what is the angular velocity of the cylinder after this collision? Is kinetic energy

conserved?



Vector Nature of Angular Quantities

The angular velocity vector points along the axis of rotation, with the direction given by the right-hand rule. If the direction of the rotation axis does not change, the angular acceleration vector points along it as well.

