

Ref: 5-1, 5-2.

e.g. Magnetic force on a rectangular current loop

Given magnetic field: $\vec{B} = \hat{x} B_0$ (constant)

$$\vec{F}_m = I \cdot \oint_C d\vec{l} \times \vec{B}$$

on (2) and (4)

$$I d\vec{l} \parallel \vec{B}$$

$$I d\vec{l} \times \vec{B} = 0.$$

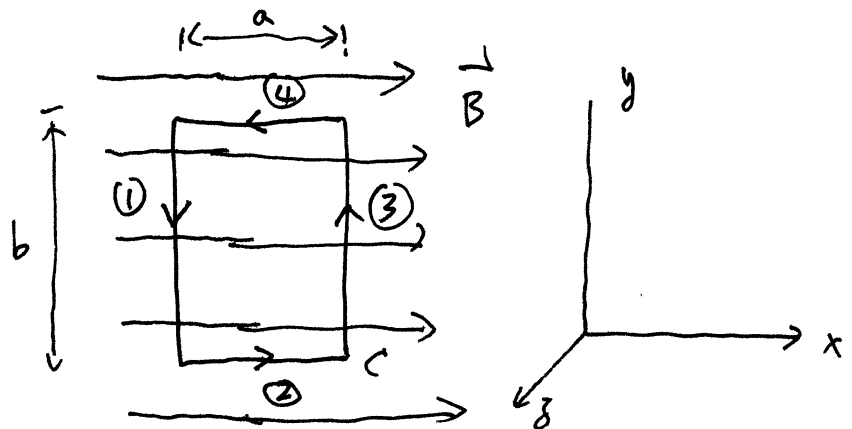
$$\therefore \vec{F}_2 = 0, \vec{F}_4 = 0 \quad (\vec{A} \times \vec{B} = 0 \text{ if } \vec{A} \parallel \vec{B}, \because \sin 0^\circ = 0)$$

$$\vec{F}_1 = I \int_{(1)} d\vec{l} \times \vec{B} = \hat{z} I b B_0.$$

$$\vec{F}_2 = -\hat{z} I b B_0.$$

$$\text{Total force: } \vec{F}_{\text{net}} = \vec{F}_1 + \vec{F}_2 + \vec{F}_3 + \vec{F}_4 = 0.$$

Centre of mass doesn't move.



But, Torque $\neq 0$ ($T = \vec{r} \times \vec{F}$)

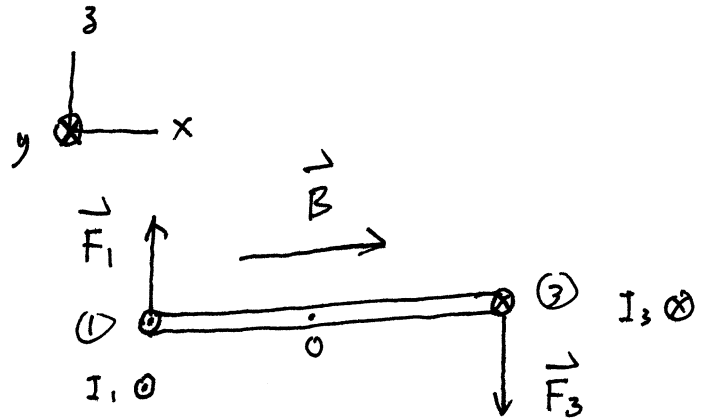
So it can rotate.

View from bottom:

Torques about O:

$$\begin{aligned}\vec{T}_1 &= (-\hat{x} \frac{a}{2}) \times \hat{z} I b B_0 \\ &= \hat{y} \frac{ab}{2} I B_0\end{aligned}$$

$$\begin{aligned}\vec{T}_3 &= (\hat{x} \frac{a}{2}) \times (-\hat{z} I b B_0) \\ &= \hat{y} \frac{ab}{2} I B_0 = \vec{T}_1\end{aligned}$$

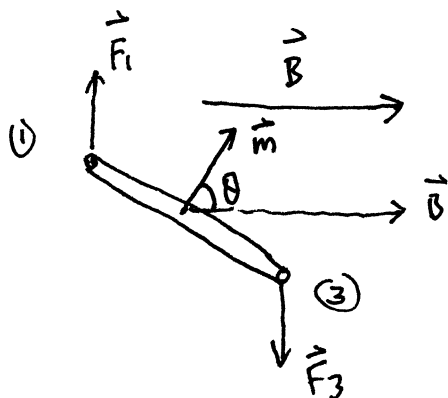


Total torque: $\vec{T} = \vec{T}_1 + \vec{T}_2 = \hat{y} \underbrace{ab \cdot I \cdot B_0}_{\text{area of loop}} = \hat{y} \vec{A} \cdot I \cdot B_0$

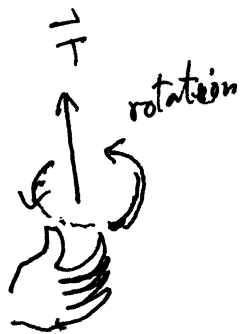
• Define magnetic moment: $\vec{m} \equiv I \cdot \vec{A}$ ($\vec{A}$ - normal)

In general: $\boxed{\vec{T} = \vec{m} \times \vec{B}} \quad (5.20)$

The loop will rotate about y-axis



axis of rotation.



$\vec{T}$: direction: $\otimes$ - meaning: RHR.
magnitude: $m B \sin \theta = T$.

• Magnetostatics

Ampere's law :
$$\oint_C \vec{H} \cdot d\vec{\ell} = I_{\text{free}}$$

$$(\vec{B} = \mu \vec{H})$$

e.g. $\vec{B}$ field of a current of a long straight line.

Along C : $\vec{B} \parallel \hat{\phi}$.
 $B = \text{constant}$.

Then:

~~$B \cdot 2\pi r I$~~

$$\oint_C \vec{B} \cdot d\vec{\ell} = \mu I$$

$$B \int_C d\ell = \mu I$$

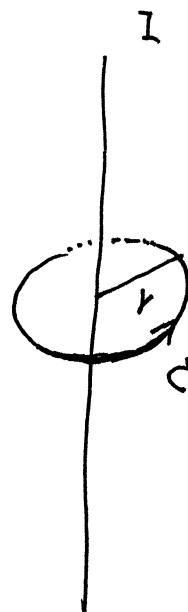
$$B 2\pi r = \mu I$$

$$\therefore \vec{B} = \hat{\phi} \frac{\mu I}{2\pi r}$$

just like Gauss' law in electrostatics.

You need symmetry !

(for $\oint_C \vec{B} \cdot d\vec{\ell}$) to make $B = \text{const.}$)

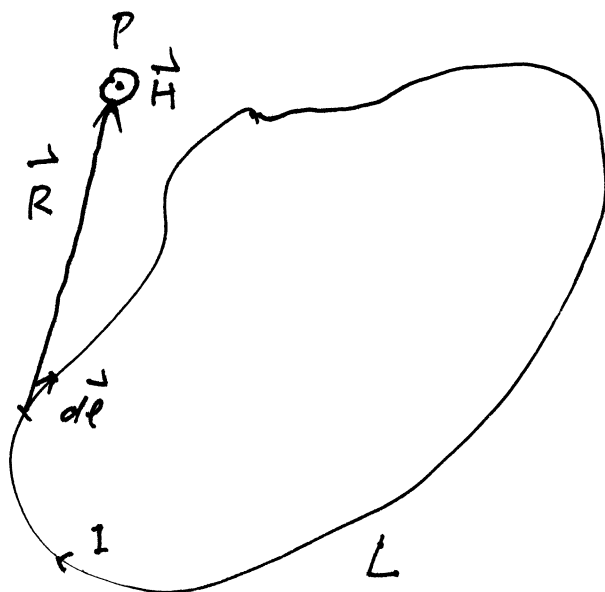


- The Biot-Savart Law.

$$d\vec{H} \text{ (due to } I d\vec{\ell} \text{)} = \frac{I}{4\pi} \frac{d\vec{\ell} \times \vec{R}}{R^2}$$

$$\vec{H} = \frac{I}{4\pi} \int_L \frac{d\vec{\ell} \times \vec{R}}{R^2}$$

(Similar to Coulomb's law
in electro statics)

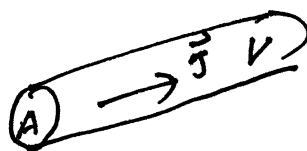


$\vec{R}$ — source to field.

- For a current distribution $\vec{J}$

$$I = \int_A \vec{J} \cdot d\vec{A}$$

Note: $\vec{J} d\vec{\ell} = \vec{J} dV$



$$\therefore \boxed{\vec{H} = \frac{1}{4\pi} \int \frac{\vec{J} \times \vec{R}}{R^2} dV} \quad (5.24)$$

$$\begin{aligned} \therefore \vec{J} dV &= (\vec{J} \cdot d\vec{\ell}) d\vec{\ell} \\ &= I d\vec{\ell} \end{aligned}$$

(Note: $\vec{J} \parallel d\vec{\ell}$)

For surface current:

$$\vec{H} = \frac{1}{4\pi} \int \frac{\vec{J}_s \times \vec{R}}{R^2} dS \quad (5.24)$$