

phys 221

Lecture 32.

Ref: 10-2, 3.

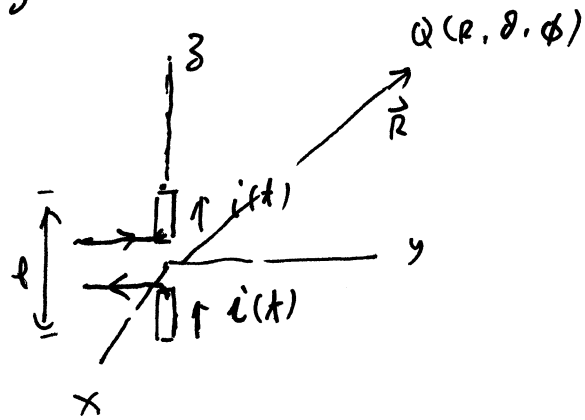
Last lecture: The short dipole (Hertzian dipole)

$$\frac{l}{\lambda} < \frac{1}{50}$$

Then, $i(t)$ ind. of z .

$$i(t) = I_0 \cos(\omega t)$$

Method: retarded vector potential.



$$\tilde{\vec{A}}(\vec{R}) = \frac{\mu_0}{4\pi} \int_{V'} \frac{\tilde{\vec{J}} e^{-jkR'}}{R'} dV' \quad \left(\text{Then } \tilde{\vec{H}}, \text{ Then } \tilde{\vec{E}} \right)$$

$$\begin{aligned} \text{Result: } \tilde{\vec{A}}(\vec{R}) &= \hat{z} \frac{\mu_0}{4\pi} I_0 l \frac{e^{-jkR}}{R} \\ &= (\hat{R} \cos \theta - \hat{\theta} \sin \theta) \frac{\mu_0 I_0 l}{4\pi} \frac{e^{-jkR}}{R} \quad (A_\phi = 0) \end{aligned}$$

spherical system.

 θ - zenith angle ϕ - azimuth angle

$$\text{Then, } \tilde{\vec{H}} = \frac{1}{\mu_0} \nabla \times \tilde{\vec{A}}$$

$$\tilde{\vec{E}} = \frac{1}{j\omega\epsilon} \nabla \times \tilde{\vec{H}}$$

$$\tilde{H}_\phi = \frac{I_0 l k^2}{4\pi} e^{-jkR} \left[\frac{j}{kR} + \frac{1}{(kR)^2} \right] \sin \theta$$

$$\tilde{E}_R = \frac{2I_0 l k^2}{4\pi} \eta_0 e^{-jkR} \left[\frac{1}{(kR)^2} - \frac{j}{(kR)^3} \right] \cos \theta$$

$$\tilde{E}_\theta = \frac{I_0 l k^2}{4\pi} \eta_0 e^{-jkR} \left[\frac{j}{kR} + \frac{1}{(kR)^2} - \frac{j}{(kR)^3} \right] \sin \theta$$

$$\eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} \approx 120\pi$$

$$(H_R = 0, H_\theta = 0, E_\phi = 0)$$

• Far field. $kR \gg 1$. $(2\pi \frac{R}{\lambda} \gg 1)$: $(\frac{1}{k^2}) (\frac{1}{k^3})$ die out

$$\tilde{E}_\theta = j \frac{I_0 l k \eta_0}{4\pi} \frac{e^{-jkR}}{R} \sin \theta \sim \frac{1}{R}$$

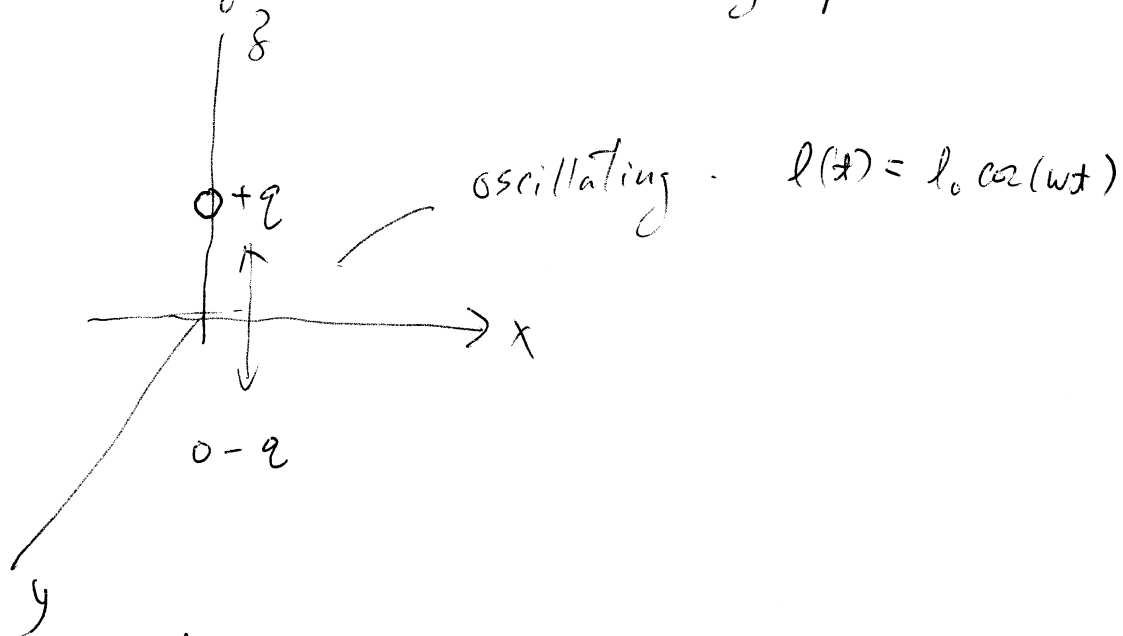
$$\tilde{H}_\phi = \frac{\tilde{E}_\theta}{\eta_0} \text{ also } \sim \frac{1}{R}$$

$$\left(\tilde{E}_R \sim \frac{1}{k^2} \eta_0 \sim 0 \right) \text{ draw}$$

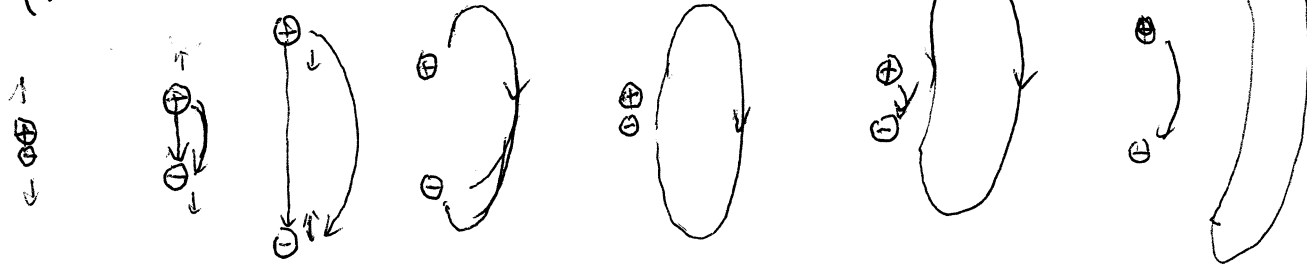
fig. 10-6. — directions of $\vec{E}$ and $\vec{H}$, $\vec{k}$ ($\vec{S}$)

fig. 10-7 — How $\vec{E}$ -field is generated and propagates

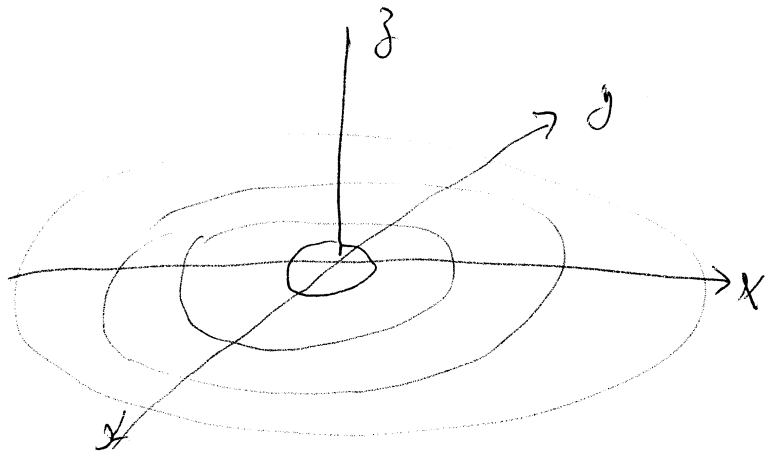
- Another way to view an oscillating dipole.



E-field.



Magnetic field: $\vec{H}$ - lines: in x - y plane



• Power density

time-average Poynting vector

$$\vec{S}_{av} = \frac{1}{2} \operatorname{Re} \left(\vec{\tilde{E}} \times \vec{\tilde{H}}^* \right) = \hat{r} S(\theta, \phi)$$

$$\begin{aligned} S(\theta, \phi) &= \frac{1}{2} \operatorname{Re} \left(\vec{\tilde{E}}_\theta \cdot \vec{\tilde{H}}_\phi^* \right) \\ &= \frac{1}{2} \operatorname{Re} \left[\frac{j I_0 \ell k \eta_0}{4\pi} \cdot \frac{e^{-jkR}}{R} \sin\theta \cdot \frac{(-j)}{\eta_0} \frac{I_0 \ell k \eta_0}{4\pi} \frac{e^{jkR}}{R} \sin\theta \right] \\ &= \frac{1}{2} \operatorname{Re} \left[\frac{I_0^2 \ell^2 k^2 \eta_0}{16\pi^2 R^2} \sin^2\theta \right] \\ &= \frac{I_0^2 \ell^2 k^2 \eta_0}{32\pi^2 R^2} \sin^2\theta \\ &= S_0 \cdot \sin^2\theta \end{aligned}$$

• Directional Pattern: $F(\theta, \phi)$

— Normalized radiation intensity

$$F(\theta, \phi) = \frac{S(R, \theta, \phi)}{S_{\max}}$$

$S_{\max}$ — Max value of $S(R, \theta, \phi)$ at a given R .

• For the Hertzian dipole

$$S_{\max} = S_0 = \frac{I_0^2 l^2 k^2 \eta_0}{32\pi^2 R^2}$$

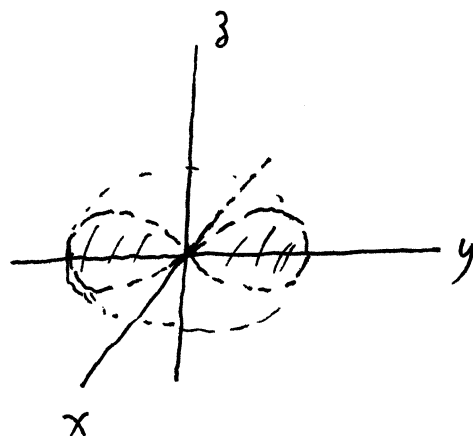
$$= \frac{15\pi I_0^2}{R^2} \left(\frac{l}{\lambda}\right)^2$$

$$F(\theta, \phi) = \sin^2 \theta$$

Figure 10-8 on page 355:

Doughnut shape.

No energy is radiated
in the z -direction.



• Elevation pattern



on the θ -plane (elevation plane) (zenith plane)

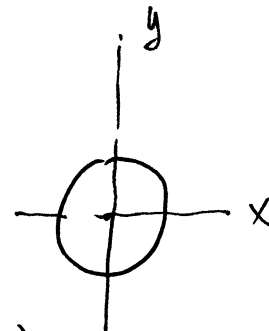
L constant ϕ .

• Azimuth pattern

on the ϕ -plane (azimuth plane)

$$\theta = 90^\circ$$

(or, just the x - y plane)

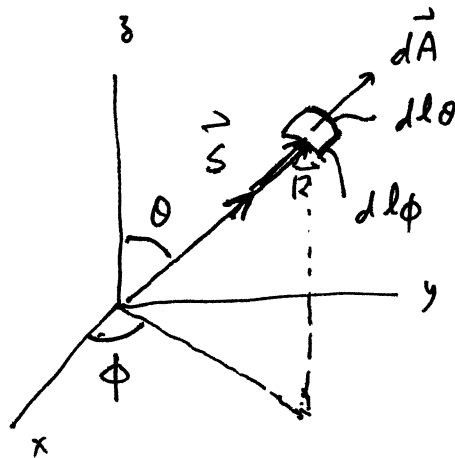


- Differential power radiated by antenna through an elemental area dA .

$$dl\theta = R d\theta$$

$$dl\phi = R \sin\theta \cdot d\phi$$

$$dA = R^2 \sin\theta \cdot d\theta \cdot d\phi$$



$$dP_{rad} = \vec{S}_{av} \cdot d\vec{A}$$

$$= (\hat{R} S) \cdot (\hat{R} dA) = S dA$$

- Solid angle associated with dA .

$$d\Omega = \frac{dA}{R^2} = \sin\theta \cdot d\theta \cdot d\phi \quad (dA = R^2 d\Omega)$$

(in sr. — steradians)

$$dP_{rad} = \underbrace{R^2 S(R, \theta, \phi)}_{\substack{\uparrow \\ \text{power per unit solid angle}}} d\Omega$$

- Total radiated power.

$$P_{rad} = R^2 \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} S(R, \theta, \phi) d\Omega = R^2 S_{max} \iint F(\theta, \phi) d\Omega$$

$$= R^2 S_{max} \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} F(\theta, \phi) \sin\theta d\theta d\phi$$

Note: Total solid angle of a ~~sp~~ sphere : $= 4\pi$.

$$\int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} d\Omega = 4\pi.$$

OR :
$$\frac{\text{surface area of sphere}}{R^2} = \frac{4\pi R^2}{R^2} = 4\pi.$$

- Antenna pattern

— described by the normalized radiation intensity.

$$F(\theta, \phi) = \frac{S(R, \theta, \phi)}{S_{\max}}.$$

How to plot $F(\theta, \phi)$ — next.