

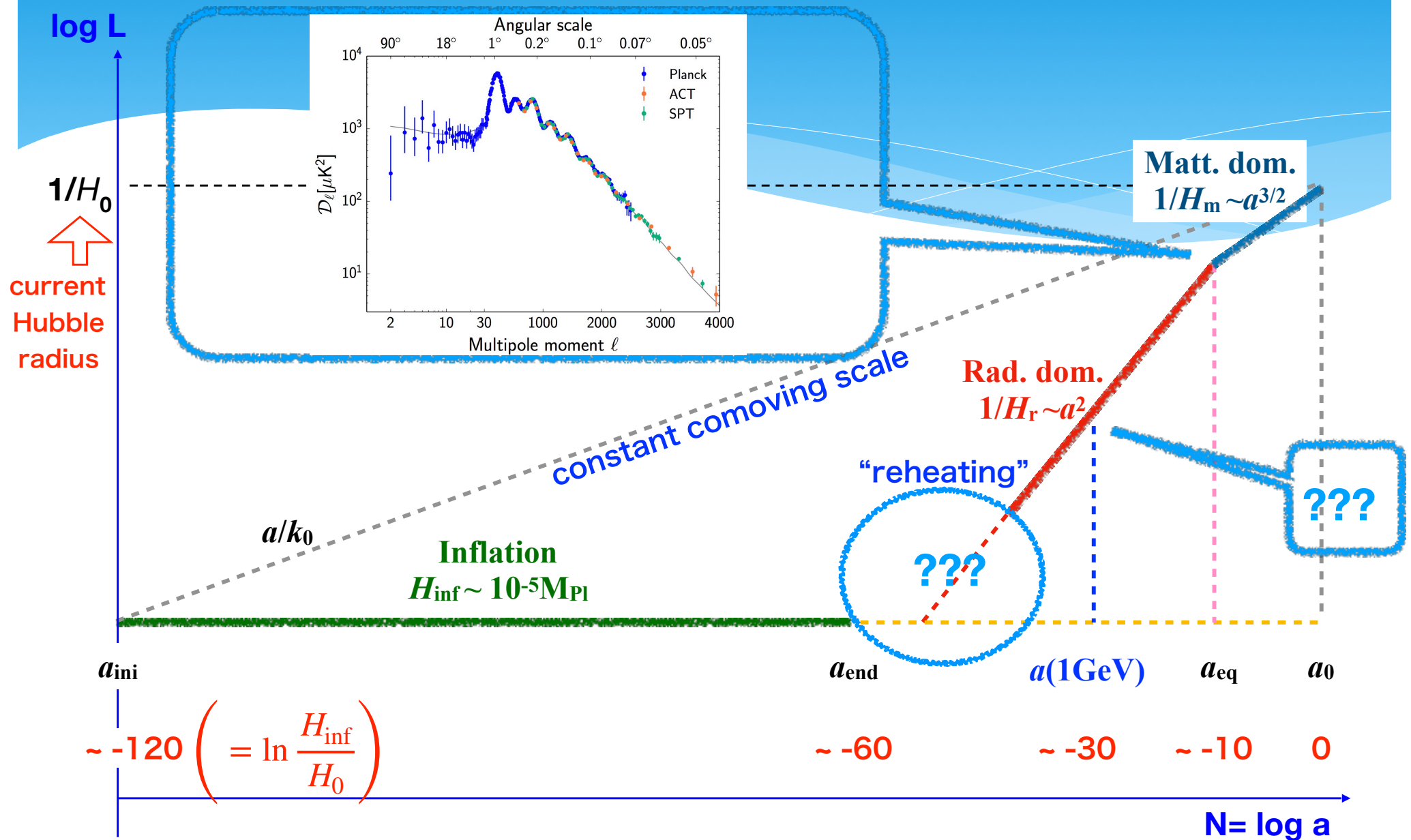
Primordial Black holes from Inflation and Gravitational Waves

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Inflation & PBH formation

cosmic spacetime diagram



curvature perturbation from inflation

- inflaton (\sim massless) vacuum fluctuations (=Gaussian)

$$\left| \langle \phi | \vec{k} \rangle \right|^2 = |\varphi_k|^2, \quad \varphi_k \sim \frac{1}{\sqrt{2\omega_k}} e^{-i\omega_k t}; \quad \omega_k = \frac{k}{a} \gg H$$

rapid expansion renders oscillations frozen at $k/a < H$
(fluctuations become “classical” on superhorizon scales)

$$\varphi_k \sim \frac{H}{\sqrt{2k^3}}; \quad \frac{k}{a} \ll H \Rightarrow \langle \delta\phi_k^2 \rangle = \left(\frac{H}{2\pi} \right)_{k/a \sim H}^2 \quad \dots \text{almost scale-invariant}$$

for $\epsilon = -\frac{\dot{H}}{H^2} \ll 1$

- curvature perturbation on comoving slices

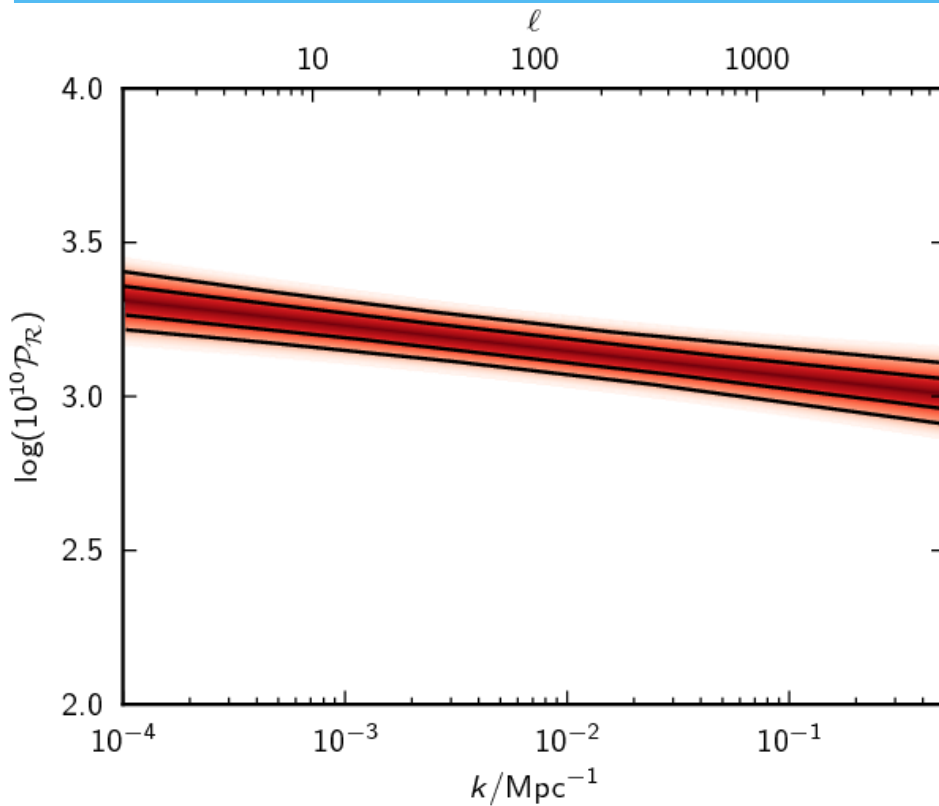
$$\mathcal{R}_c = -\frac{H}{\dot{\phi}} \delta\phi \quad \dots \text{conserved on superhorizon scale}$$

for single-field slow-roll models

(almost scale-invariant if $\dot{\phi}$ is also slowly varying)

observational constraint on inflation

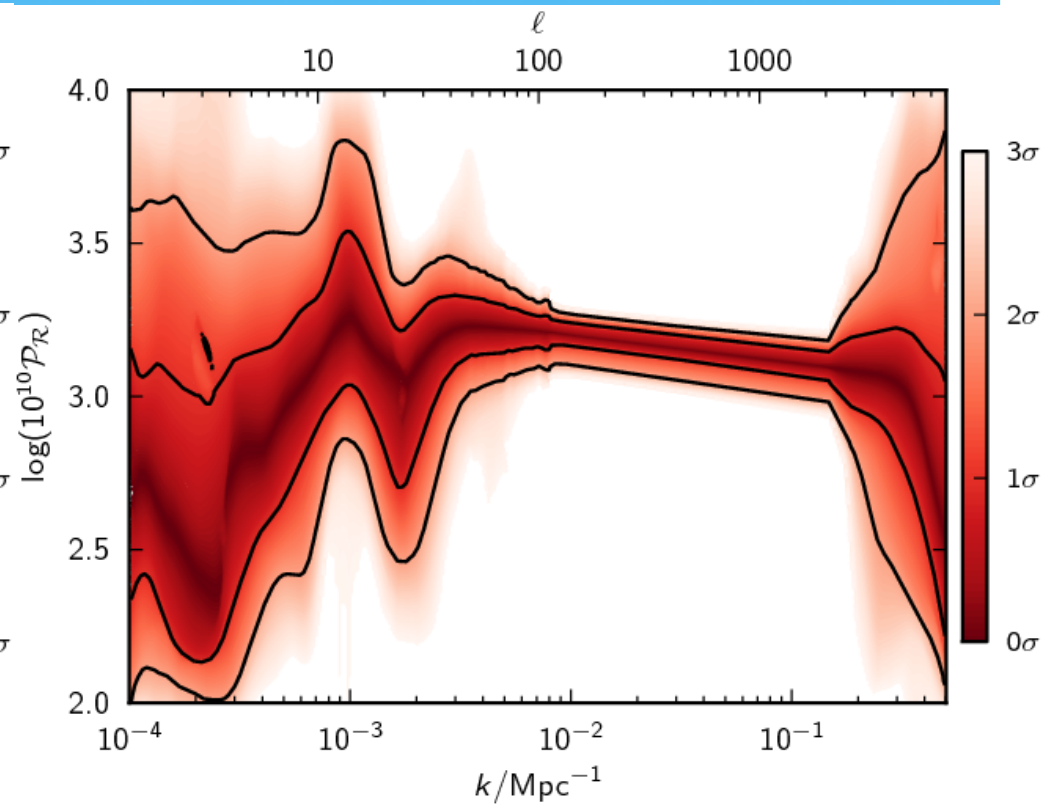
Planck 2015 results XX



power-law

$$\mathcal{P}_{\mathcal{R}} = A(k/k_*)^{n_s-1}$$

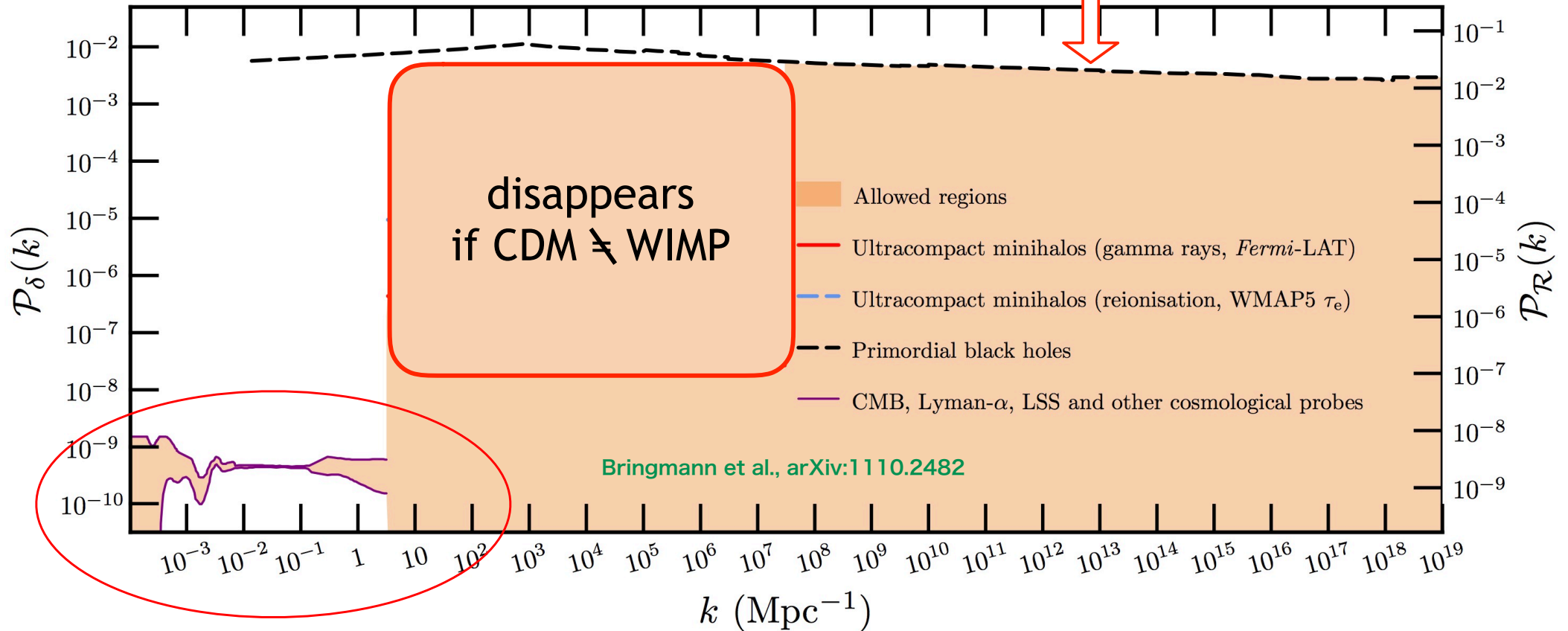
$n_s \approx 0.968 \dots$ almost scale-invariant



piece-wise continuous
(9 segments)

observational constraint on inflation

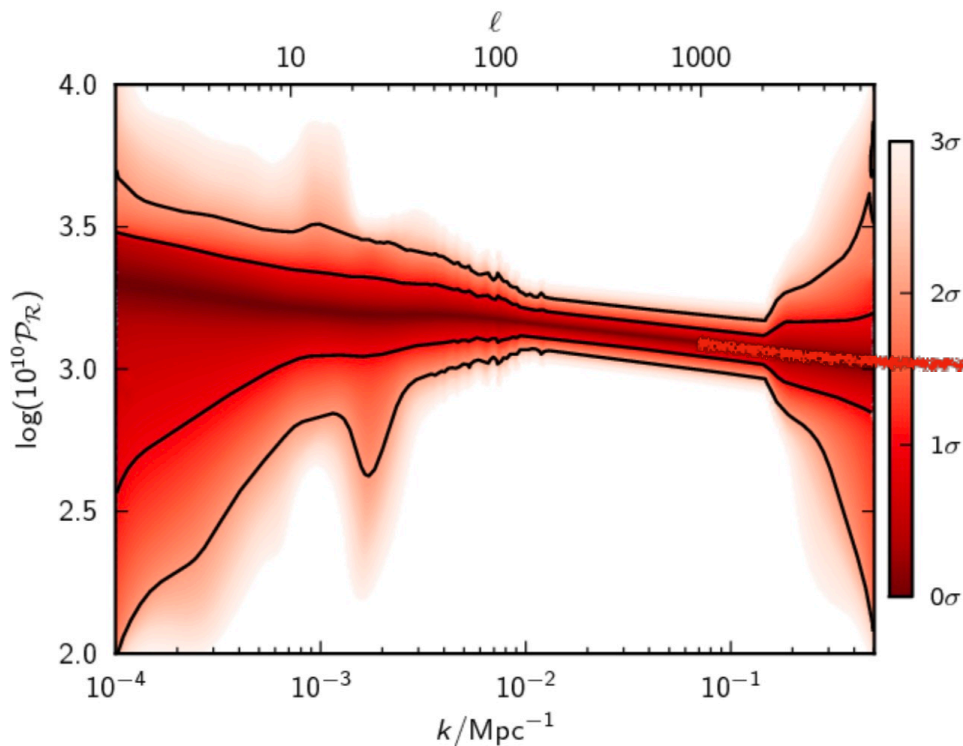
constraints on small scales are from BHs



Planck constraint

There are some constraints on small scales, but quite weak.

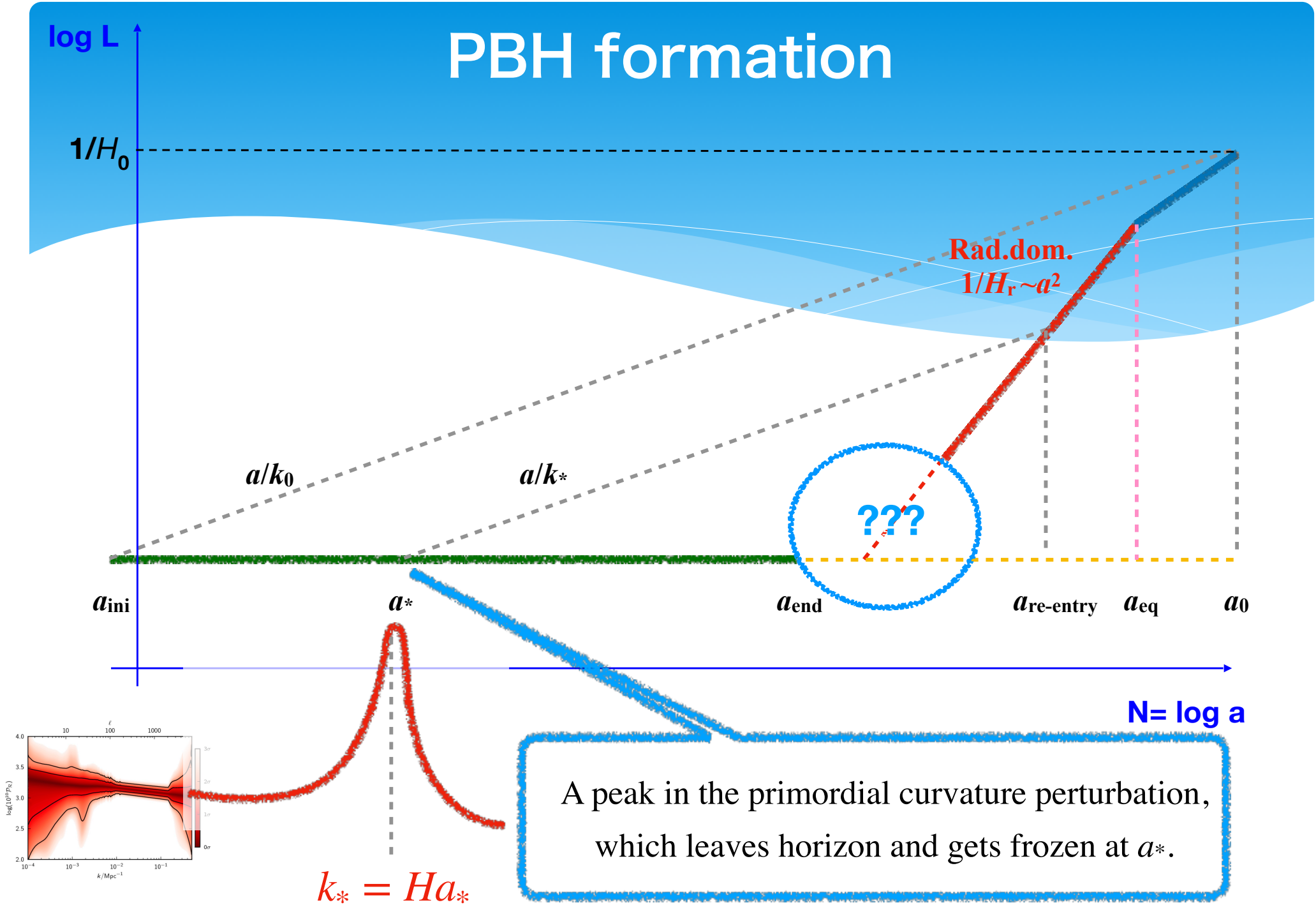
Bayesian reconstruction of the primordial power spectrum with $l < 2300$. (Planck 2015)



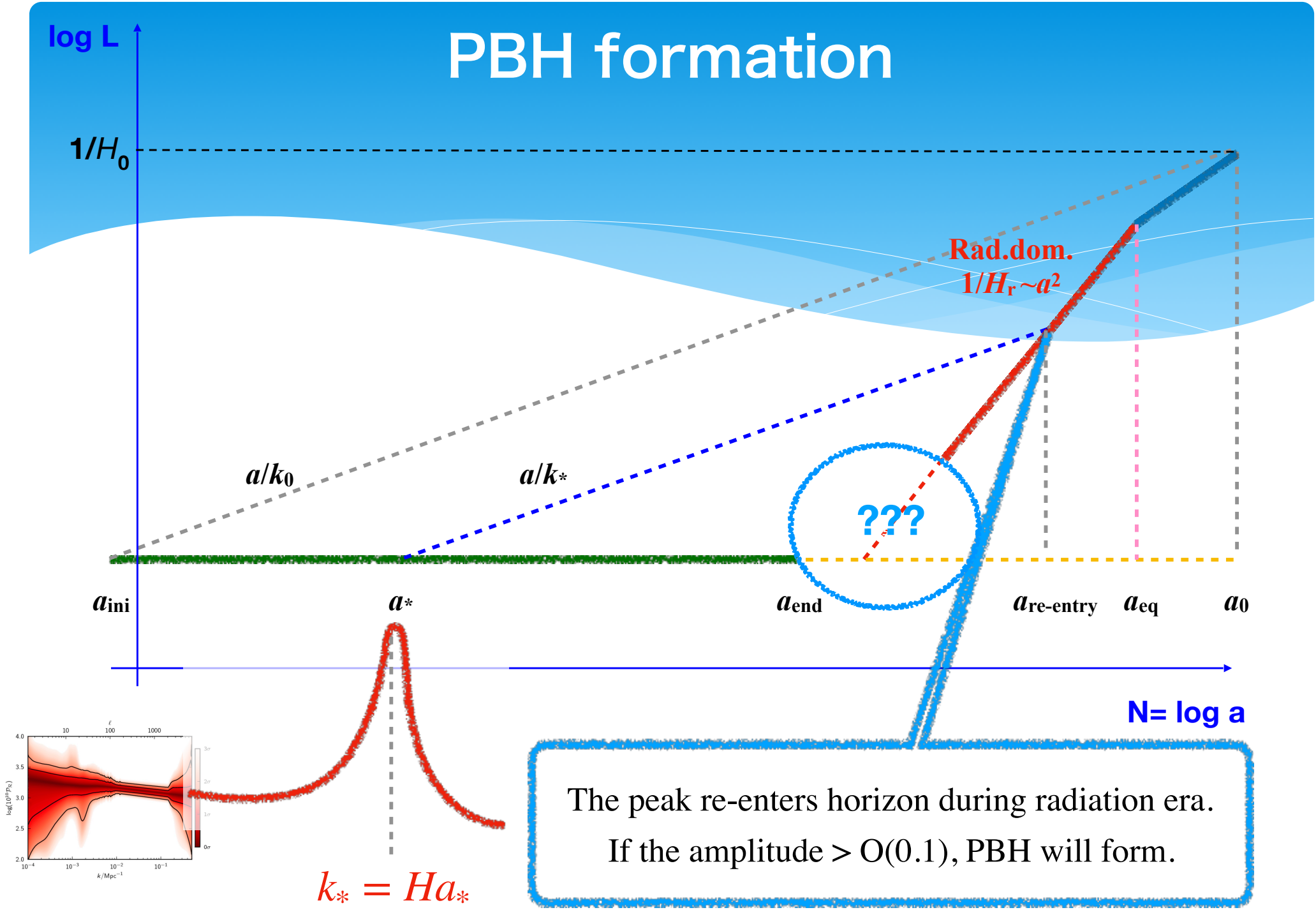
No resolution to say anything precise about higher k .

?

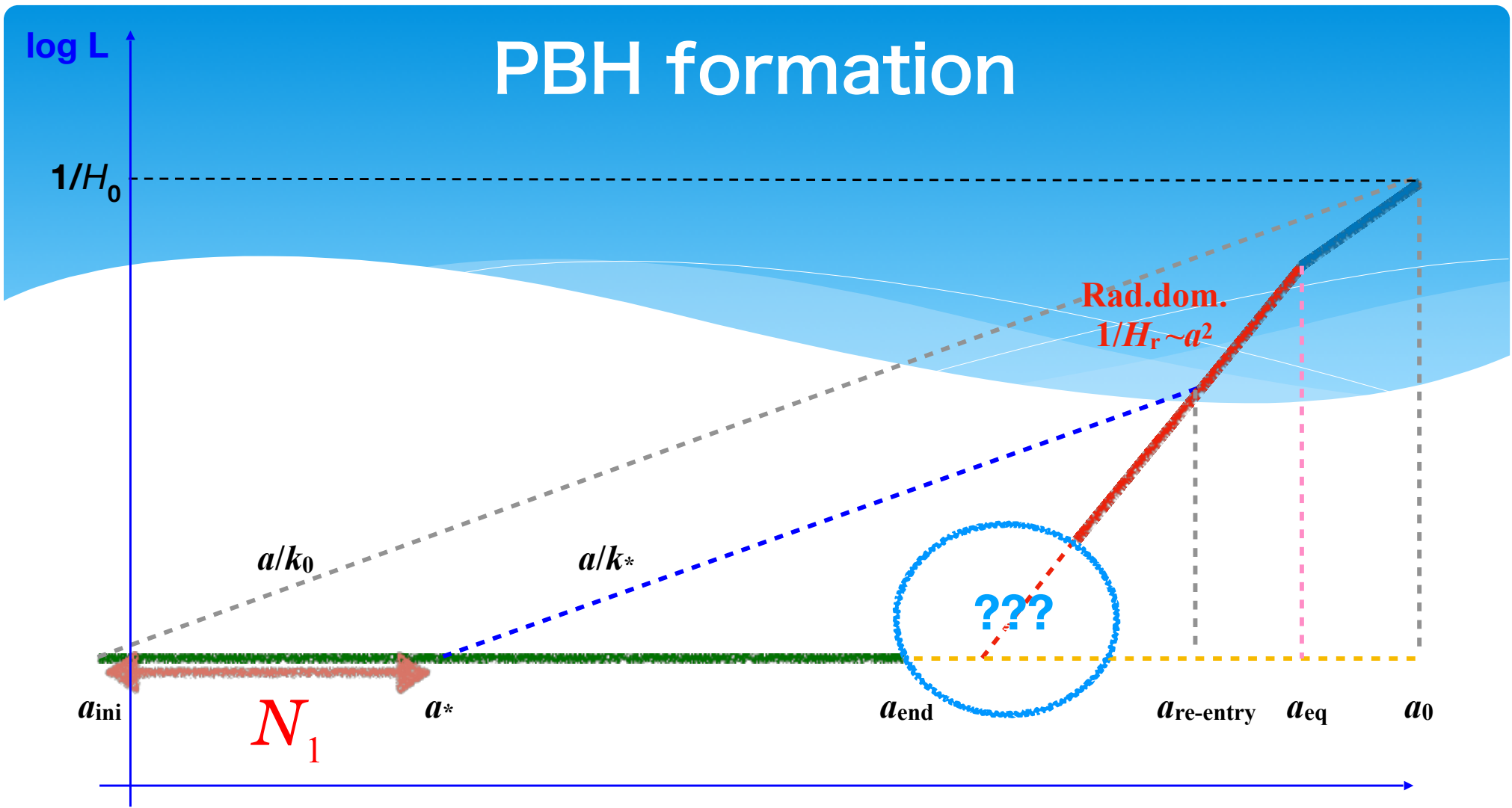
PBH formation



PBH formation



PBH formation



$N = \log a$

PBH mass: $M_{PBH} = \gamma M_H \sim \frac{M_{Pl}^2}{H} = 10^{58} M_{Pl} e^{-2N_1} = M_{Pl} 10^{58-0.87N_1}$

Inverse relation: $N_1 = 44.4 - \frac{1}{2} \ln \left(\frac{M_{PBH}}{10^{16} \text{g}} \right)$

PBH mass scale does NOT depend on the reheating physics

Primordial Black Holes

What are Primordial BHs?

- **PBH = BH formed before recombination epoch** (ie at $z \gg 1000$)
conventionally during radiation-dominated era
- **Hubble size region with $\delta\rho/\rho=O(1)$** collapses to form BH
Carr (1975),
- Such a large perturbation may be **produced by inflation**
Carr & Lidsey (1991), ...
- PBHs may dominate **Dark Matter**.
Ivanov, Naselsky & Novikov (1994), ...
- **Supermassive BHs** ($M \gtrsim 10^6 M_\odot$) may originate from PBHs.
Bean & Magueijo (2002), ...
-

Curvature perturbation to PBH

➤ gradient expansion/separate universe approach

$$6H^2(t, \mathbf{x}) + R^{(3)}(t, \mathbf{x}) = 16\pi G\rho(t, \mathbf{x}) + \dots \quad \text{Hamiltonian constraint (Friedmann eq.)}$$

$$\Rightarrow R^{(3)} \approx -\frac{4}{a^2} \nabla^2 \mathcal{R}_c \approx \frac{8\pi G}{3} \delta\rho_c \quad \Rightarrow \quad \frac{\delta\rho_c}{\rho} \approx \mathcal{R}_c \quad \text{at} \quad \frac{k^2}{a^2} = H^2$$

$$R^{(3)} \simeq 0 \quad \left\langle \begin{array}{c} R^{(3)} \sim H^2 \\ \longleftrightarrow \\ H^{-1} = a/k \end{array} \right\rangle$$

➤ If $R^{(3)} \sim H^2$ ($\Leftrightarrow \delta\rho_c / \rho \sim 1$) collapses to form BH

Young, Byrnes & MS '14

$$M_{\text{PBH}} \sim \rho H^{-3} \sim 10^5 M_{\odot} \left(\frac{t}{1\text{s}} \right) \sim 20 M_{\odot} \left(\frac{k}{1\text{pc}^{-1}} \right)^{-2}$$

➤ Spins of PBHs are expected to be very small

examples

hybrid-type inflation

Garcia-Bellido, Linde & Wands '96, ...

\mathcal{R}_C grows near the saddle point
non-Gauss may become large

Abolhasani, Firouzjahi & MS '11, ...
 Pattison et al. 1707.00537

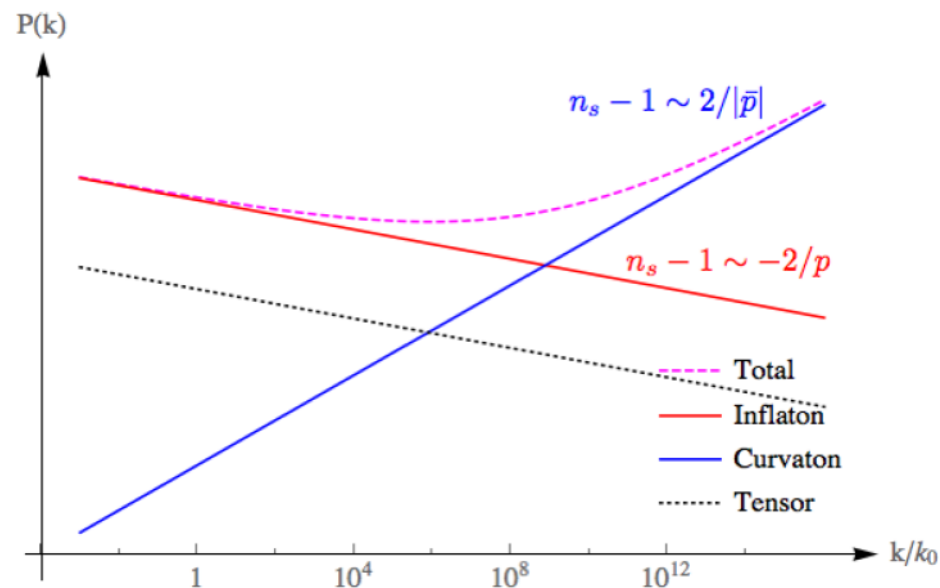
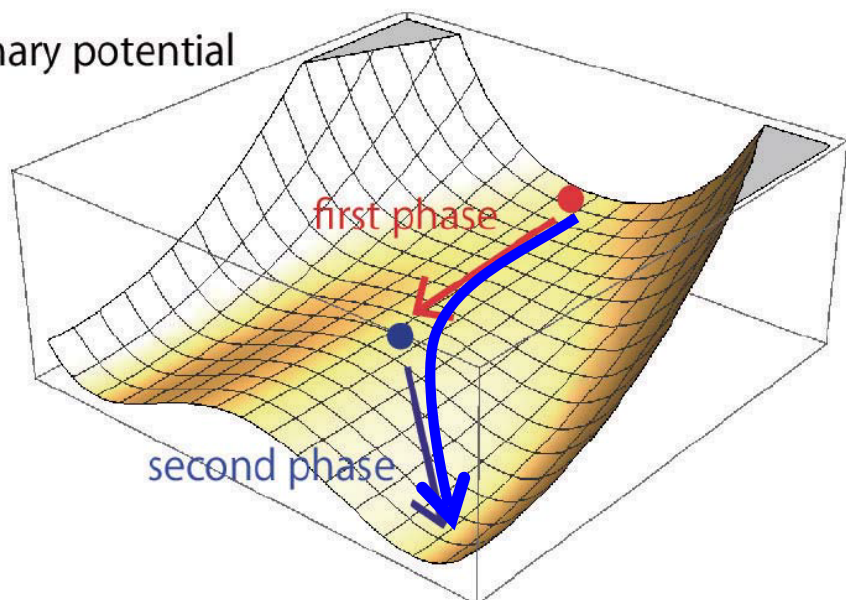
non-minimal curvaton

Domenech & MS '16

$$L = -\frac{1}{2} f(\phi) g^{\mu\nu} \partial_\mu \chi \partial_\nu \chi$$

$$-\frac{1}{2} h(\phi) m^2 \chi^2$$

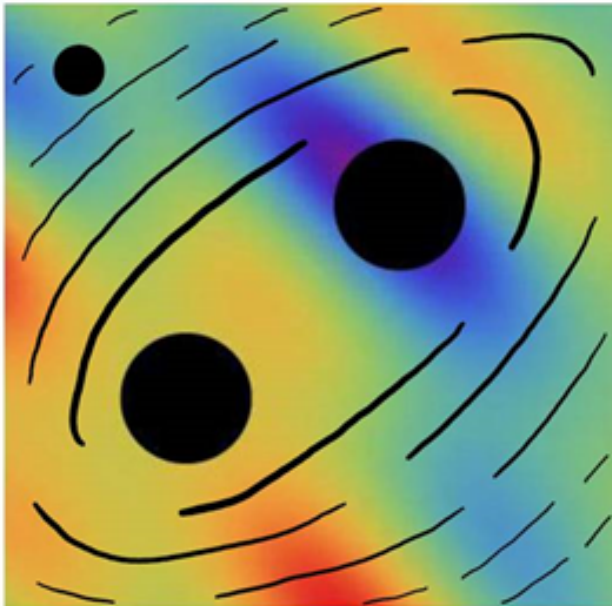
inflationary potential



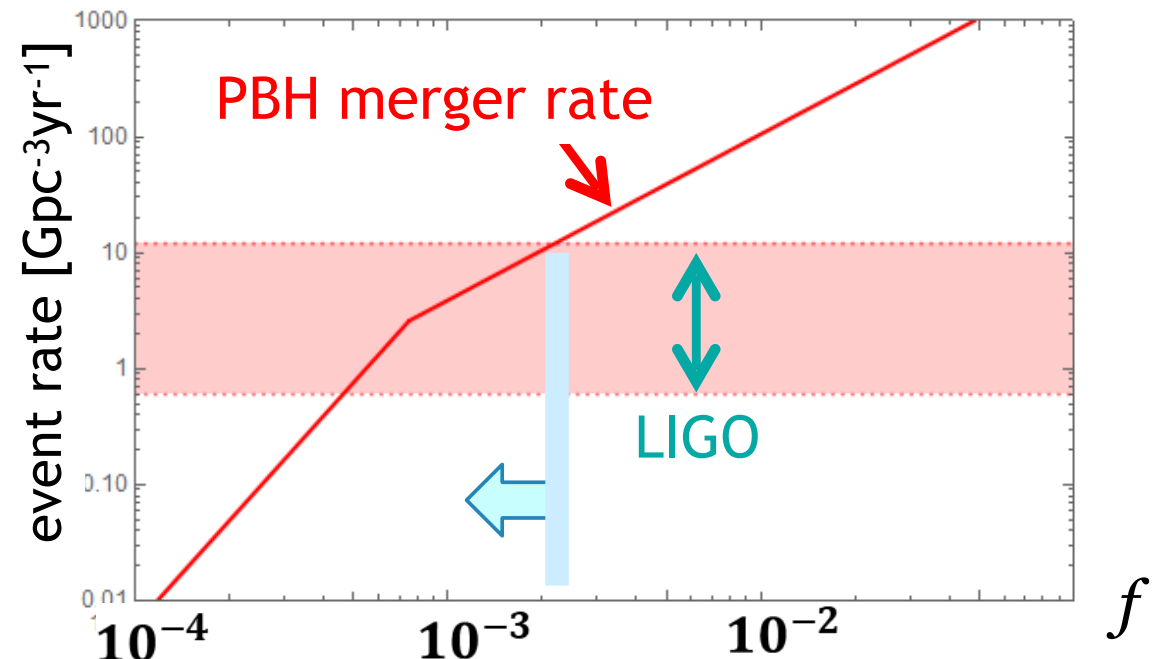
LIGO BHs = PBHs?

MS, Suyama, Tanaka & Yokoyama '16

$$M_{PBH} \simeq 20 \left(\frac{k}{\text{kpc}^{-1}} \right)^{-2} M_{\odot} \simeq 20 \left(\frac{100 \text{MeV}}{T} \right)^2 M_{\odot}$$



3-body interaction leads to formation of BH binaries

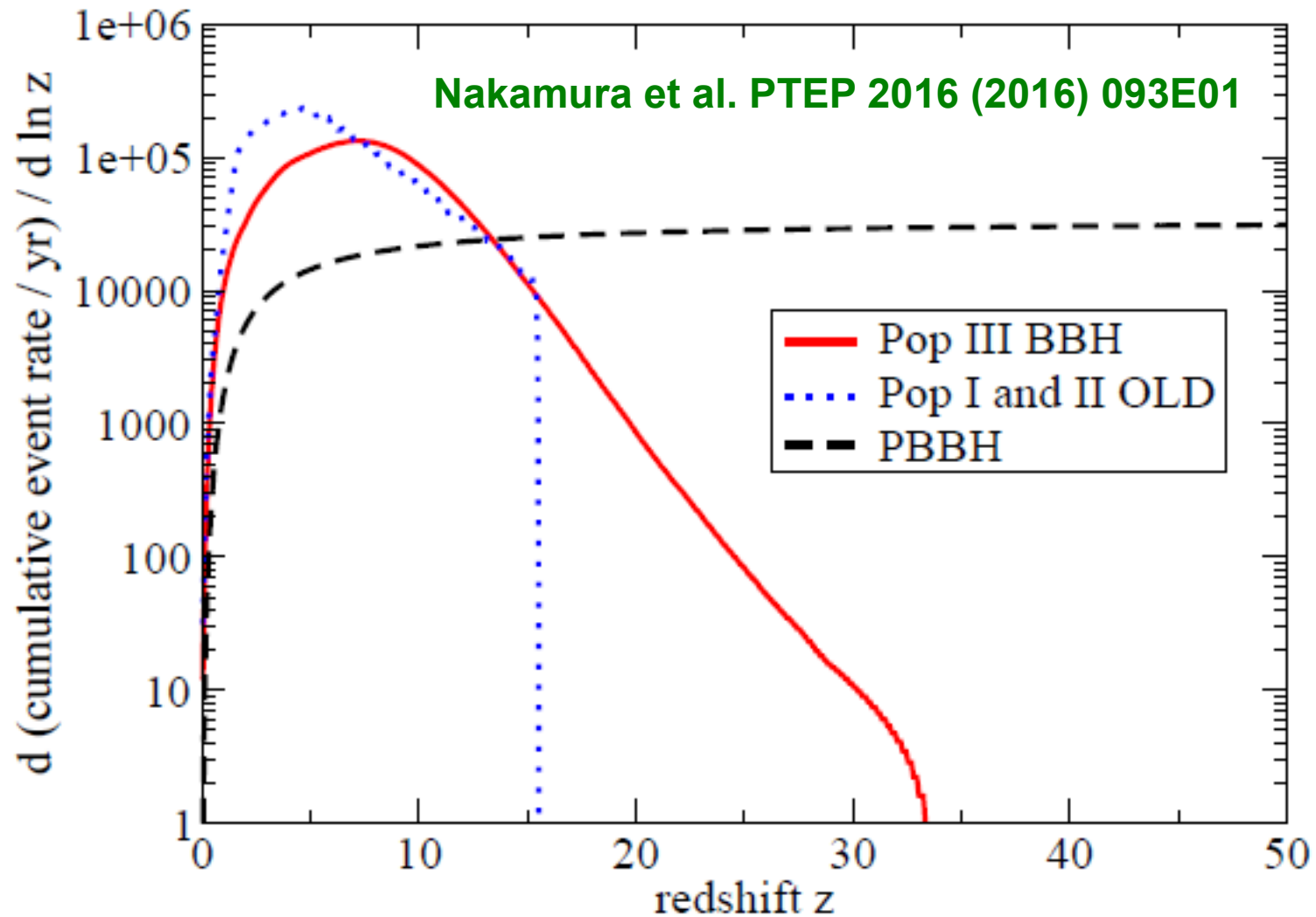


f = fraction of PBH in dark matter

tightest constraint at $M \sim 10M_{\odot}$

(cf. Ali-Haïmoud et al., 1709.06576)

testing PBH hypothesis



testing PBH hypothesis 2

Kocsis, Suyama, Tanaka, Yokoyama, arXiv:1709.09007

BBH Merger Rate at time t : mass function

$$\mathcal{R}(m_1, m_2, t) = \frac{n_{\text{BH}}}{2} f(m_1) f(m_2) P_{\text{intr}}(m_1, m_2, t)$$

intrinsic probability

$$P_{\text{intr}}(m_1, m_2, t) \propto g(m_1) g(m_2) m_t^\alpha : m_t = m_1 + m_2$$

$$\iff \alpha(m_1, m_2, t) \equiv -m_t^2 \frac{\partial^2}{\partial m_1 \partial m_2} \ln \mathcal{R}(m_1, m_2, t)$$

- PBH binary scenario

$$\frac{36}{37} < \alpha < \frac{22}{21}$$

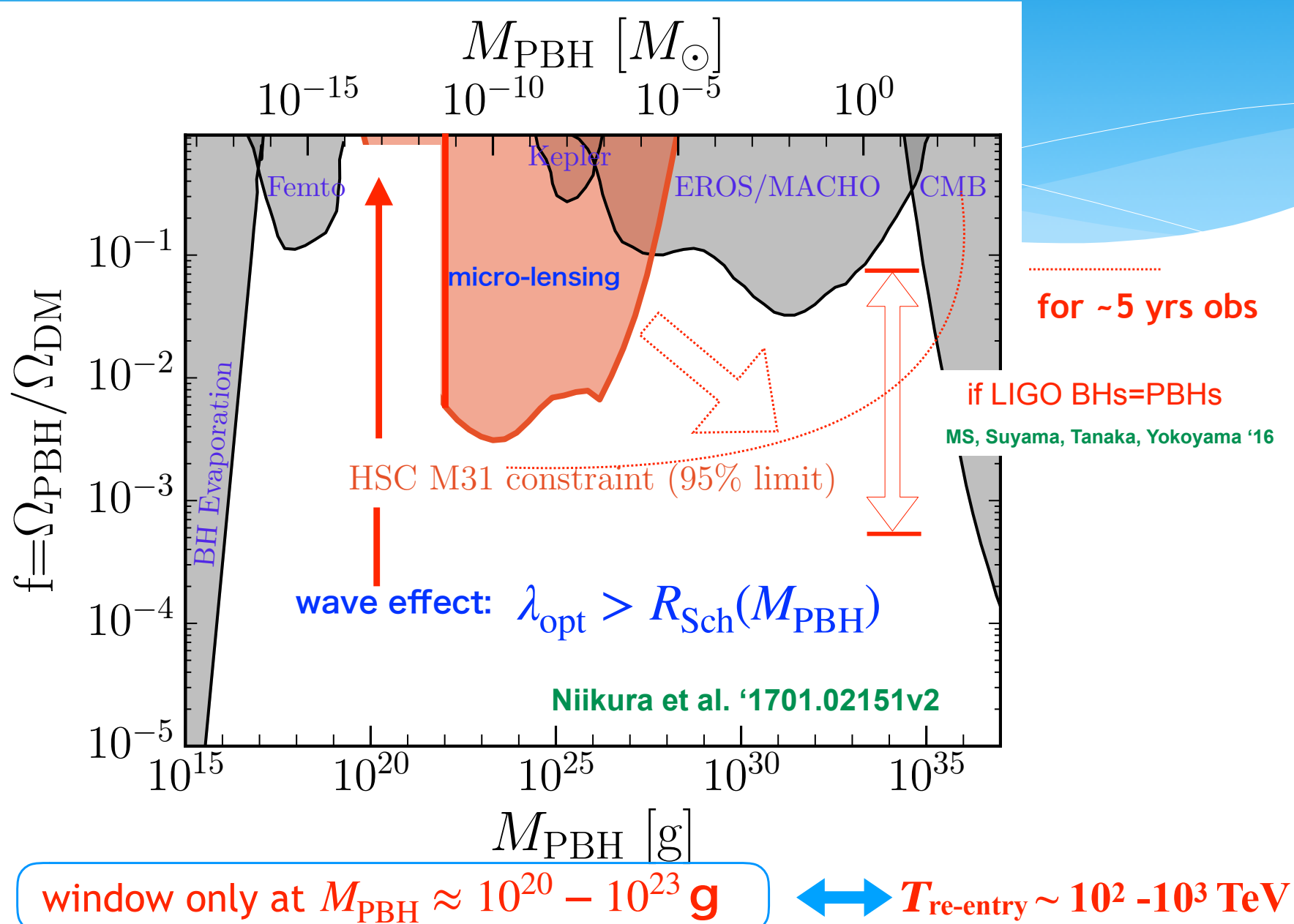
- Dynamical formation in dense stellar systems

$$\alpha \approx 4$$

clearly distinguishable!

O'Leary et al (2016)

PBH constraints: recent update

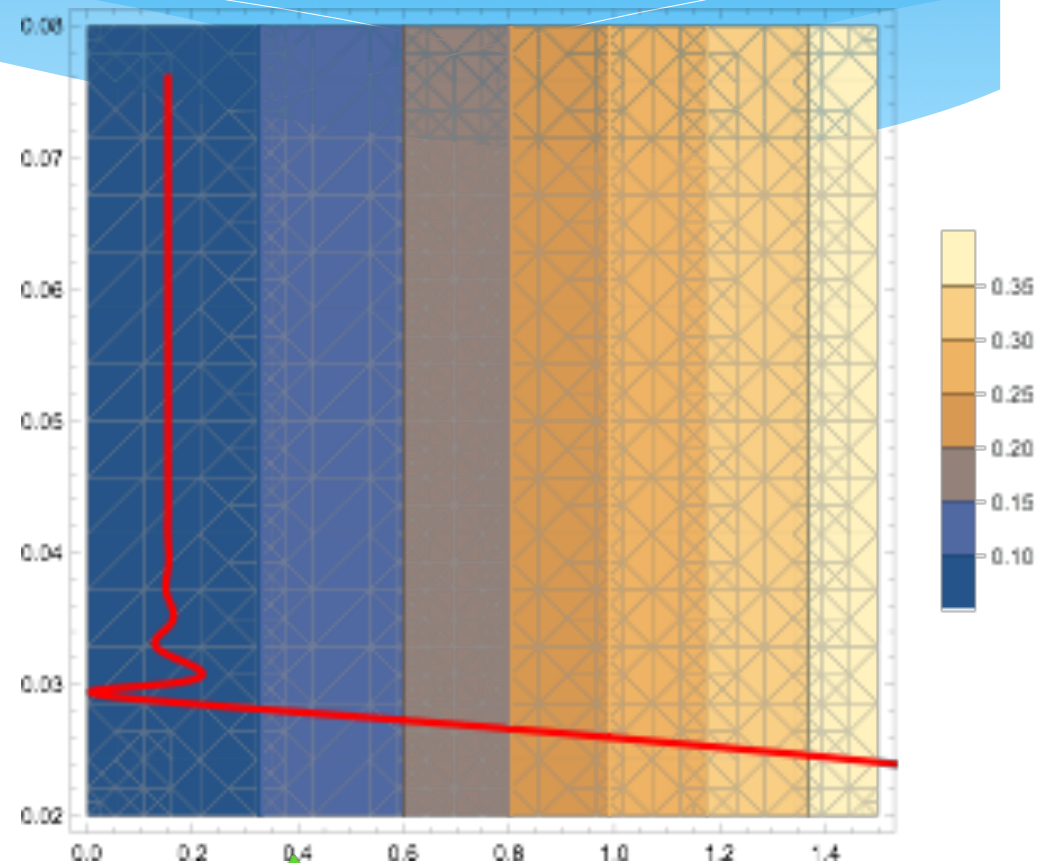
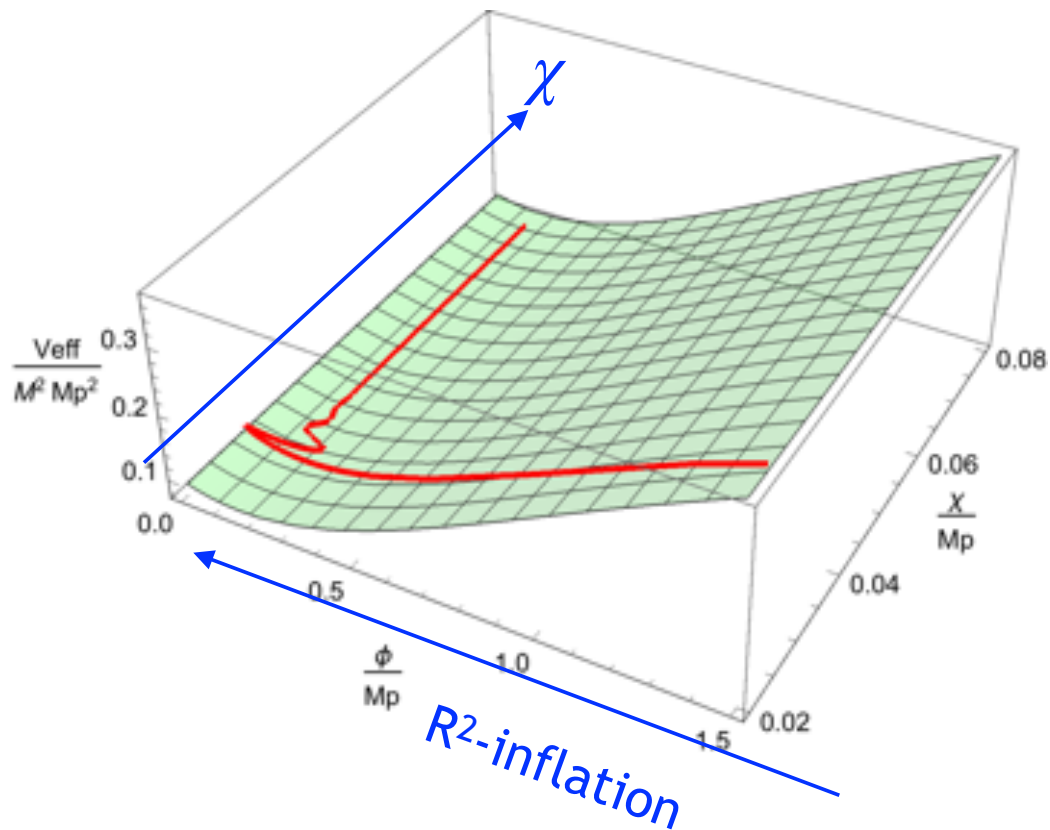


PBHs = CDM?

monochromatic PBH production

Pi, Zhang, Huang & MS '18

$$\mathcal{L} = R + \frac{R^2}{6M^2} - \frac{1}{2}(\partial\chi)^2 - V(\chi)$$

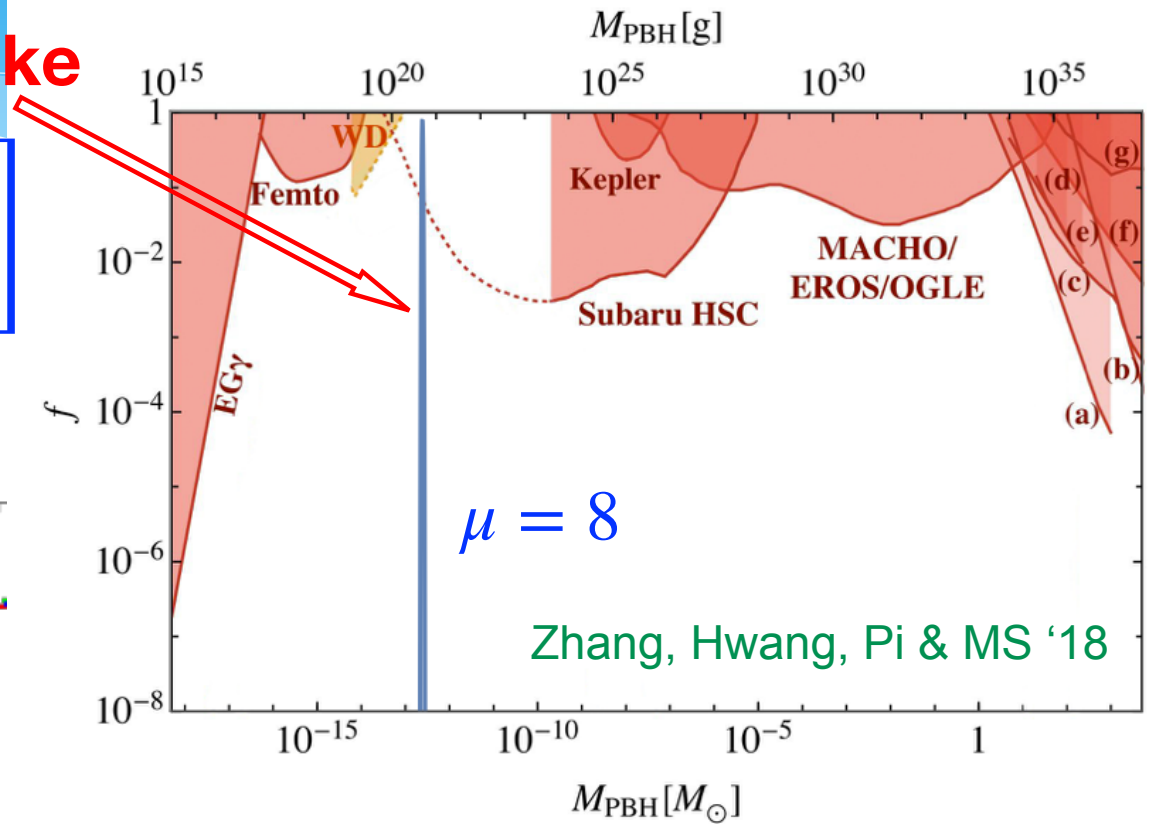


End of the 1st stage of inflation

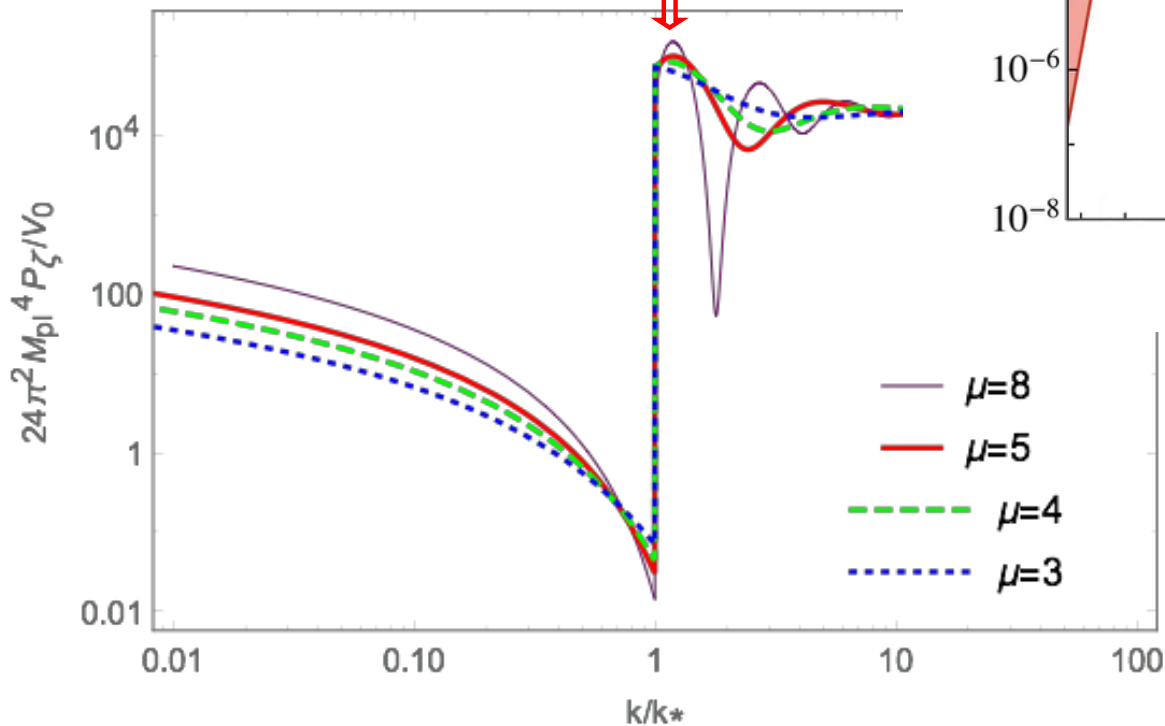
sharp peak in $P(k) \rightarrow$ spike in f

$$f \equiv \frac{\Omega_{\text{PBH}}}{\Omega_{\text{CDM}}} \propto \exp \left[-\frac{1}{\mathcal{P}(k)} \right]$$

spike



sharp peak



Zhang, Hwang, Pi & MS '18

$$\mu^2 \approx \frac{H_{2\text{nd}}^2}{H_{1\text{st}}^2}$$

2nd order GWs

Saito & Yokoyama '09, Alabidi et al. '13, ...

- non-negligible PBH formation means $\mathcal{P}_S(k) \gtrsim 10^{-2.5}$

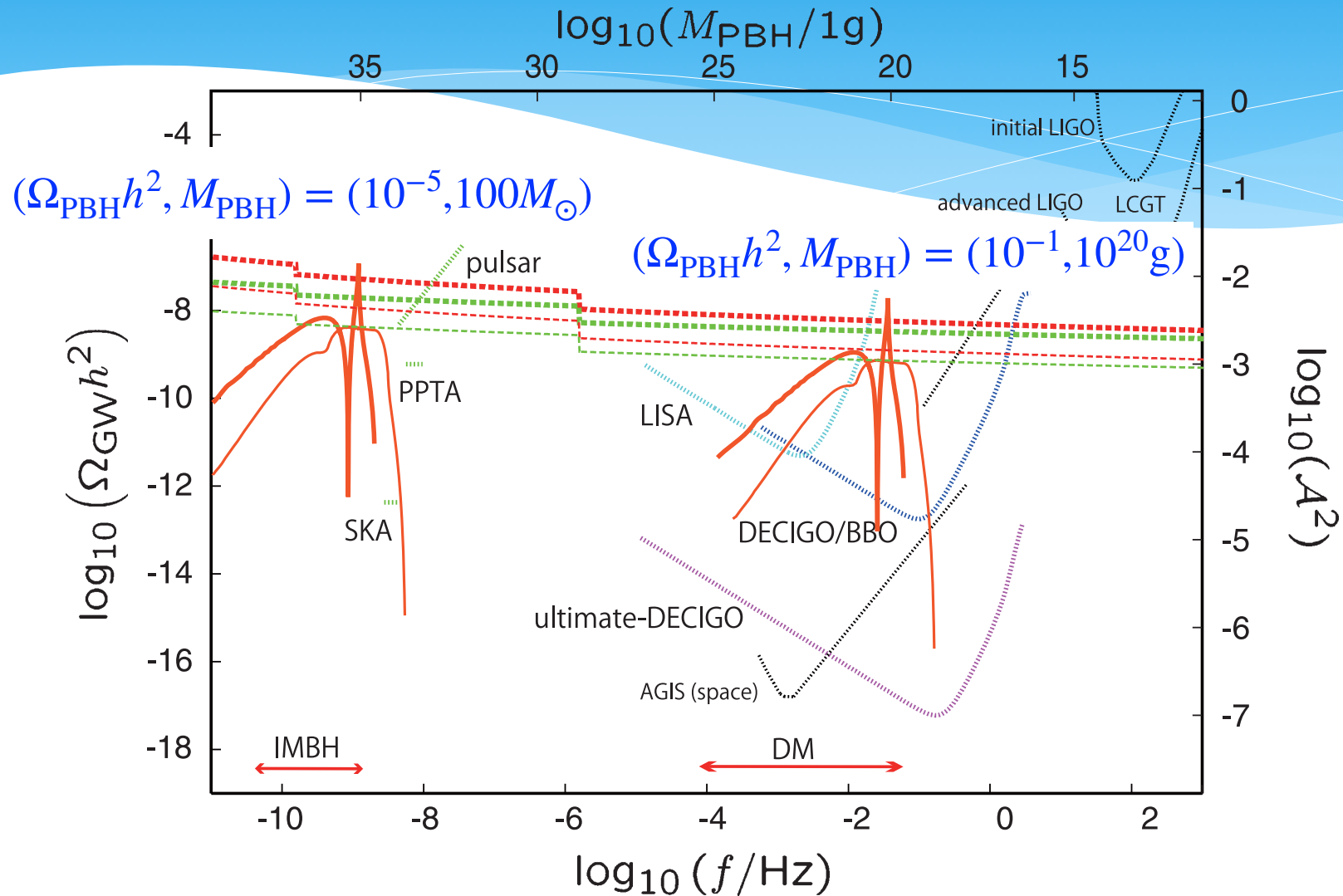
$$\ddot{h}_{ij} + 3H\dot{h}_{ij} - a^{-2}\Delta h_{ij} = S_{ij}$$
$$S_{ij} = \frac{1}{a^2}\partial_i\mathcal{R}_c\partial_j\mathcal{R}_c + \dots \sim \frac{k^2}{a^2}\mathcal{P}_S(k)$$

- GWs are produced with amplitude:

$$h_{ij} \sim \frac{k^2}{a^2 H^2} \mathcal{P}_S(k) \sim \mathcal{P}_S(k) \Rightarrow \mathcal{P}_{GW}(k) \sim (\mathcal{P}_S(k))^2$$

2nd order GWs could dominate GW background at $f > 10^{-10}$ Hz ($k > 10^4$ Mpc $^{-1}$)

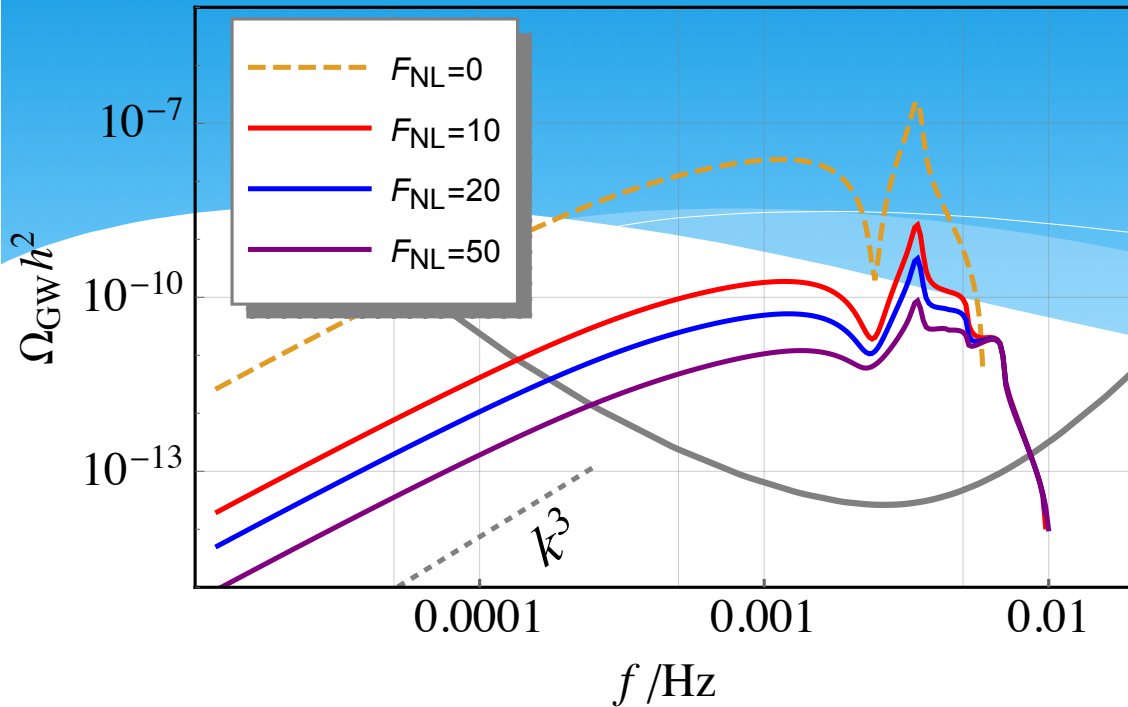
GWs test PBH scenarios!



Saito & Yokoyama, arXiv:0912.5317

Non-Gaussianity Effect?

RG Cai, S Pi & MS, arXiv:1810.11000



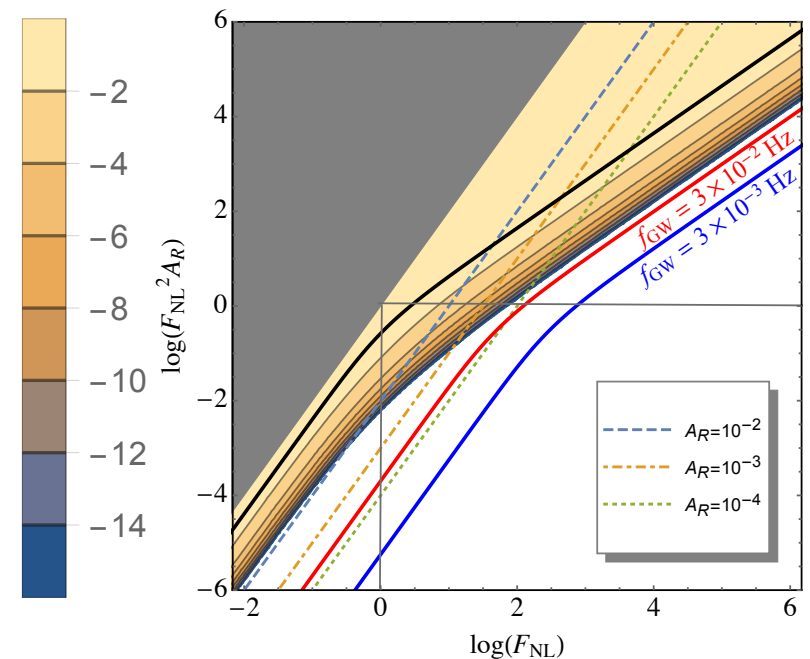
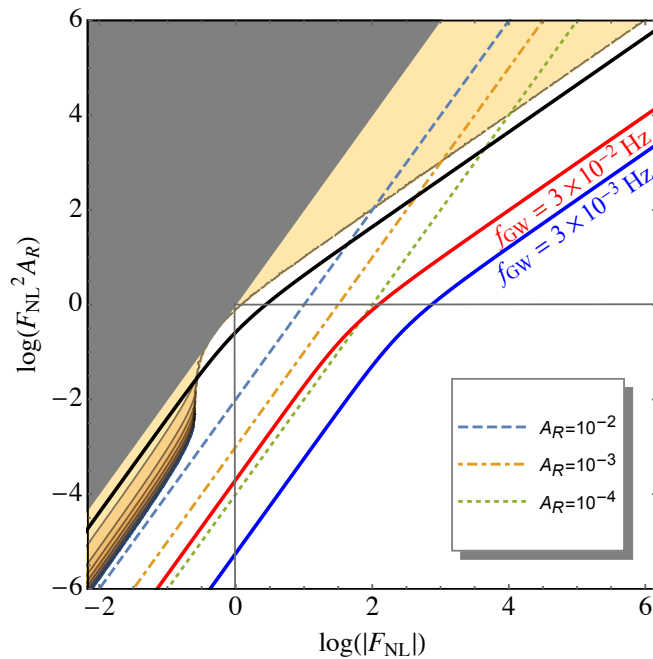
$$\mathcal{R}(x) = \mathcal{R}_g(x)$$

$$+ F_{\text{NL}} \left[\mathcal{R}_g^2(x) - \langle \mathcal{R}_g^2(x) \rangle \right]$$

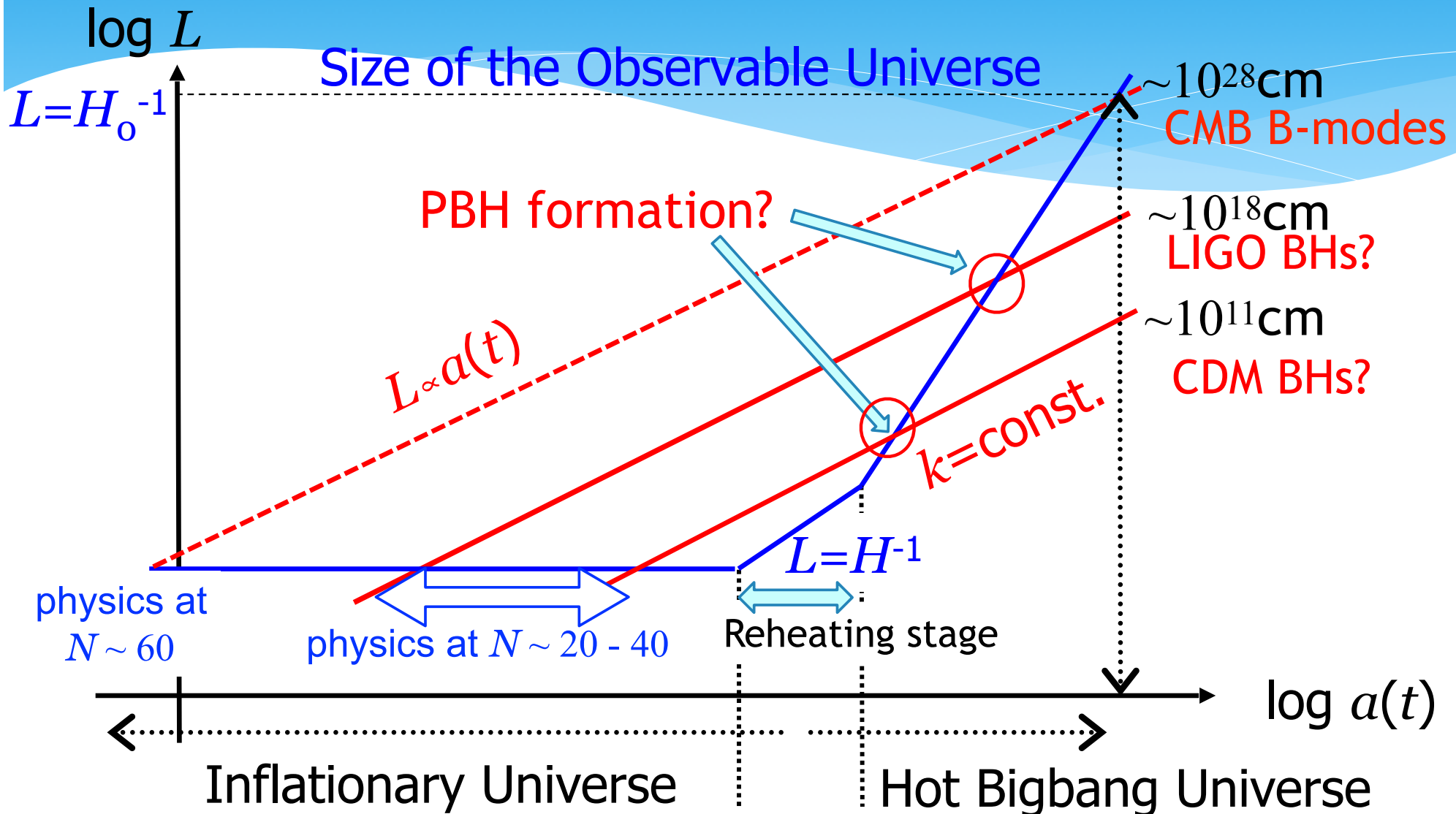
$$M_{\text{PBH}} = 10^{22} \text{ g}$$

PBH constraints:

If PBHs=CDM
LISA **must detect**
induced GWs!



probing inflation by GW astronomy



Summary

- * **Inflation** has become the **standard model** of the Universe.
further tests are needed to confirm inflation.
- * Inflation can produce **large curvature perturbation** on **small scales**.
PBHs are virtually the only probe on very small scales.
- * **LIGO BHs** may be **primordial**.
future GW detectors(+G lensing) will prove/disprove the scenario.
- * **CDM** can be dominated by PBHs of $M \sim 10^{20} \text{g}$.
secondary GWs may be detected by future GW detectors.
- * **Multi-frequency GW** astronomy/astrophysics is arriving!

GWs will be **an essential tool** for proving/
falsifying PBH scenarios