

Physics. 101.

Lecture 22.

Wed. Feb. 26, 2003.

Topics Covered.

Simple Harmonic Motion.

— A vibrating block.

— Simple pendulum.

Ref: 13-1, 2, 3, 4, 5, 6

- Simple Harmonic Motion

characteristic equation: $a = -\omega^2 x$.

acceleration $\propto (-)$ displacement.

$$x = A \cdot \cos(\omega t + \varphi)$$

$$v = -\omega A \sin(\omega t + \varphi)$$

$$a = -\omega^2 A \cdot \cos(\omega t + \varphi)$$

3 constant parameters describing the oscillation:

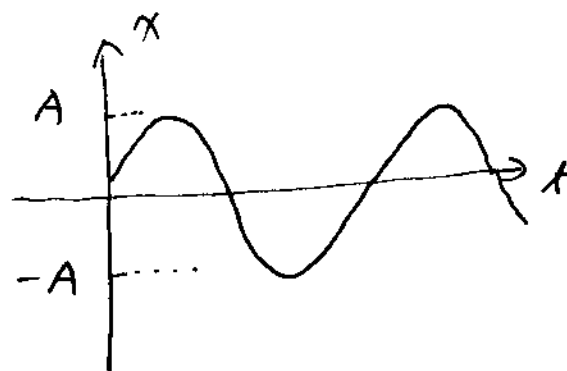
A, ω, φ .

- A — amplitude

$$x_{\max} = A$$

$$v_{\max} = \omega A$$

$$a_{\max} = \omega^2 A$$



- ω — angular frequency (How quickly it oscillates)

period: $T = \frac{2\pi}{\omega}$, $\therefore \omega = \frac{2\pi}{T}$.

frequency: $f = \frac{1}{T} = \frac{\omega}{2\pi}$

- φ — initial phase.

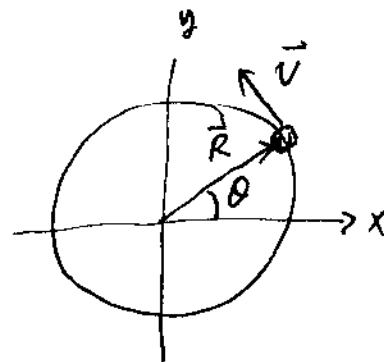
How and when the oscillation starts.

Text book: special case $\varphi = 0$

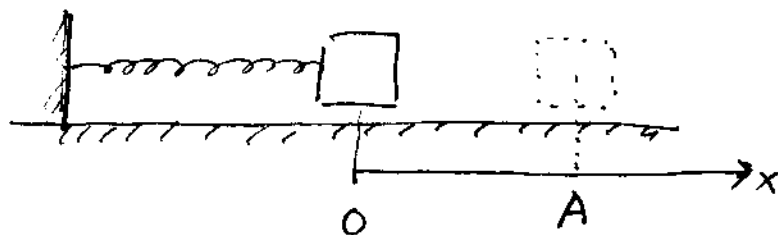
Dr. Boal's Notes: special case $\varphi = -\frac{\pi}{2}$.

Note: It's a good idea

to think about the
reference circle (~~rotating~~
ping-pong on a rotating wheel).



e.g: If we pull the block to $x = A = x_{\max}$,
Then release it at $t = 0$.



i.e., when $t = 0$, $x = A$, $v = 0$

Then $A = A \cdot \cos(0 + \varphi) \Rightarrow \cos \varphi = 1$, $\varphi = 0$.

$$\therefore x = A \cos(\omega t)$$

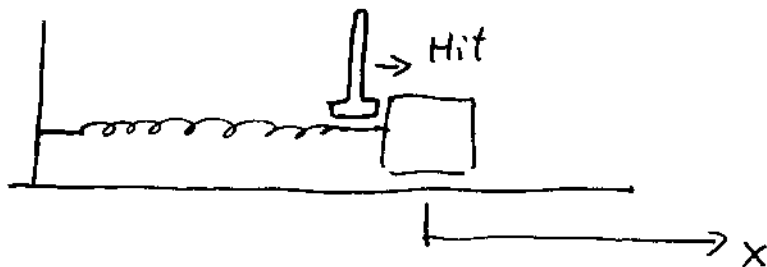
$$v = -\omega A \sin(\omega t)$$

$$a = -\omega^2 A \cdot \cos(\omega t)$$

(Text book case)

another example,

If we give the block an initial velocity $v = \omega A = v_{\max}$



by hitting it to the right at $x = 0$ and $t = 0$.

i.e., when $t = 0$, $x = 0$, $v = \omega A = v_{\max}$.

$$\left. \begin{aligned} \text{Then, } 0 &= A \cdot \cos(0 + \varphi) \\ \omega A &= -\omega A \cdot \sin(0 + \varphi) \end{aligned} \right\} \Rightarrow \begin{aligned} \cos \varphi &= 0, \varphi = \pm \frac{\pi}{2} \\ \sin \varphi &= -1, \varphi = -\frac{\pi}{2} \end{aligned}$$

$$\therefore \left\{ \begin{aligned} x &= A \cdot \cos\left(\omega t - \frac{\pi}{2}\right) = A \cdot \sin(\omega t) \\ v &= -\omega A \cdot \sin\left(\omega t - \frac{\pi}{2}\right) = \omega A \cos(\omega t) \\ a &= -\omega^2 A \cdot \cos\left(\omega t - \frac{\pi}{2}\right) = -\omega^2 A \cdot \sin(\omega t) \end{aligned} \right. \quad \left. \begin{array}{l} \text{Dr. Boal's} \\ \text{case.} \end{array} \right.$$

Note: In general, ϕ can have any value.

$\therefore$ A and ϕ can be determined by initial conditions. (How and when the oscillation is started).

However, ω is determined by dynamics.
e.g. { Force, mass, etc }.

e.g.: A vibrating Block.

Hooke's Law:

$$F = -kx = ma$$



$$\therefore a = -\frac{k}{m}x$$

compare to: $a = -\omega^2 x$ [characteristic eq. of SHM]

$$\therefore \omega^2 = \frac{k}{m}, \quad \omega = \sqrt{\frac{k}{m}}$$

k — Spring constant . (How strong the spring is)

The greater the k , the faster it oscillates .

m — mass .

The larger the m , the slower it oscillates .
(harder to accelerate)

• Mechanical Energy . at any time .

$$E = U + K .$$

$$= \frac{1}{2} k x^2 + \frac{1}{2} m v^2 .$$

~~≡~~ If there is no friction , E will be conserved .
i.e , $E = \text{const}$.

Math :
$$E = \frac{1}{2} k [A \cdot \cos(\omega t + \phi)]^2 + \frac{1}{2} m [-\omega A \cdot \sin(\omega t + \phi)]^2$$

$$= \frac{1}{2} k A^2 \cdot \cos^2(\quad) + \frac{1}{2} m \omega^2 A^2 \cdot \sin^2(\quad)$$

$$= \frac{1}{2} k A^2 \cdot \cos^2(\quad) + \frac{1}{2} k A^2 \cdot \sin^2(\quad)$$

$$= \frac{1}{2} k A^2 [\cos^2(\quad) + \sin^2(\quad)]$$

$$= \frac{1}{2} k A^2 = \text{const} .$$

Note: $\omega^2 = \frac{k}{m}$
 $m \omega^2 = k$ $\Rightarrow$

• Simple Pendulum

Can be considered as
rotation about A :

$$\tau_{\text{net}} = -mgL \cdot \sin \theta$$

$$I = mL^2$$

Since $\tau = I\alpha$. (α - angular acc.)

$$-mgL \cdot \sin \theta = mL^2 \alpha$$

$$\alpha = -\frac{g}{L} \cdot \sin \theta$$

When θ is small, $\sin \theta \approx \theta$ (if θ is in radians)

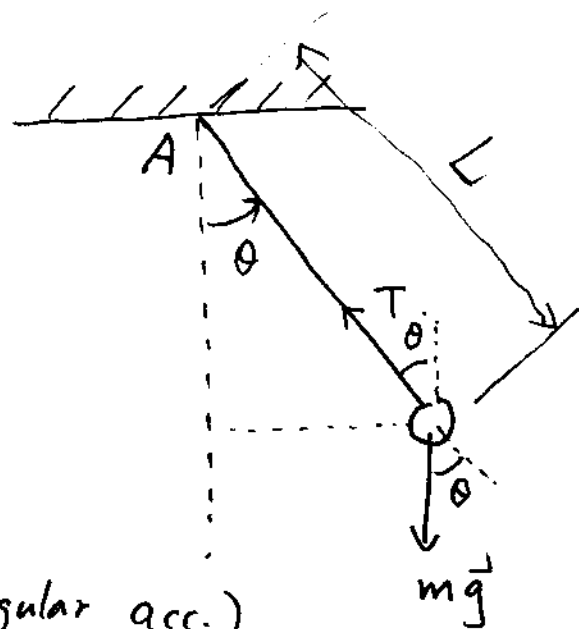
$$\therefore \alpha \approx -\frac{g}{L} \cdot \theta$$

i.e. angular acceleration $\propto (-)$ angular displacement.

compared to . $a = -\omega^2 x$ (characteristic eq. for SHM)

We have : $\omega^2 = \frac{g}{L}$. $\omega = \sqrt{\frac{g}{L}}$

$$\theta = A \cdot \cos(\omega t + \varphi) \quad \text{--- SHM .}$$



Meaning: • The larger the g , the faster it oscillates.
 $\therefore$ restoring force is larger.

$\therefore$ restoring force is larger.

- ω is independent of m .

A simple pendulum can be used to measure the value of g .