Spontaneous Change

For a system in thermal equilibrium with its surroundings,

 $dS \geqslant \frac{\delta q}{T}$ Clausius inequality

At constant volume:

$$dq_V = dU$$
 no work
 $TdS - dU \ge 0$
 $("dS)_{U,V} \ge 0$ isolated system
 $(dU)_{S,V} \le 0$

At constant pressure:

$$dq_{P} = dH \qquad PV \text{ work only}$$
$$TdS - dH \ge 0$$
$$(dH)_{S,P} \le 0$$

For convenience, define:

A = U - TSdA = dU - TdS - SdTG = H - TSdG = dH - TdS - SdT

Then the conditions for spontaneous change become:

$$(dA)_{T,V} \leq 0$$
 $(dG)_{T,P} \leq 0$

Helmholtz Energy A

Helmholtz energy; Helmholtz free energy; Helmholtz function; Maximum work function A for Arbeit

For spontaneous change at constant T and V

 $dA = dU - TdS \leqslant 0$

Note that it is the *total* function *A* that tends to a minimum; this is *not* the same as minimizing *U* and maximizing *S*.

Maximum Work

Combine $dU = \delta q + \delta w$ and $TdS \ge \delta q$

 $\Rightarrow dU \leq TdS + \delta w$

equality for reversible change

Word done by

the system $(-\delta w) \leq TdS - dU$

 $(-\delta w)_{\rm max} = TdS - \delta q_{\rm rev} - \delta w_{\rm rev} = -\delta w_{\rm rev}$

A system does maximum work when it is operating reversibly.

But
$$(dA)_T = dU - TdS = dU - \delta q_{rev} = \delta w_{rev}$$

Therefore, for macroscopic changes

$$-w_{\rm max} = -\Delta A = T\Delta S - \Delta U$$
 constant T

(-*w*) can be more or less than ΔU according to the sign of ΔS . For $\Delta S > 0$, heat flows into the system to fuel the extra work.

Gibbs Energy G

Gibbs free energy, Gibbs function

Very important in chemistry since it tells whether a particular reaction can proceed at a given T and P.

For spontaneous change,

$$(dG)_{T,P} \leqslant 0$$
 $\Delta G = G_{\text{products}} - G_{\text{reactants}} \leqslant 0$

 $\Delta G_{T,P}$ for reactions can be calculated from tabulated data $\Delta G^{\circ}(T) = \Delta H^{\circ}(T) - T\Delta S^{\circ}(T)$

| | | · · · · · · · · · · · · · · · · · · · | |
|------------------|-------------------|---------------------------------------|------------------------------|
| If ΔH is | and ΔS is | then ΔG | the reaction proceeds |
| -ve | +ve | < 0 | at all temperatures |
| +ve | -ve | > 0 | at no temperatures |
| -ve | -ve | | if $T < \Delta H / \Delta S$ |
| +ve | +e | | if $T > \Delta H / \Delta S$ |
| | | | |

Maximum Work:

$$(dG)_{T} = dH - TdS$$

$$(dG)_{T,P} = dU + PdV - TdS$$

$$= \delta q_{rev} + \delta w_{rev} + PdV - TdS$$

$$= \delta w_{rev} + PdV$$

$$= \delta w_{max} (non-PV)$$

$$\Delta w_{max} (non-PV) = \Delta G \qquad \text{constant } T, P$$

Basic Thermodynamic Relations 1

| Laws | $dU = \delta q - PdV$ | (1) |
|-------------|-------------------------------|-----|
| | $dS = \delta q_{\rm rev} / T$ | (2) |
| Definitions | H = U + PV | (3) |
| | A = U - TS | (4) |
| | G = H - TS | (5) |

Fundamental Equations

| dU = TdS - PdV | (6) |
|-----------------|-----|
| dH = TdS + VdP | (7) |
| dA = -SdT - PdV | (8) |
| dG = -SdT + VdP | (9) |

Partial Differentials

$$\begin{pmatrix} \frac{\partial U}{\partial S} \end{pmatrix}_{V} = T \qquad \begin{pmatrix} \frac{\partial U}{\partial V} \end{pmatrix}_{S} = -P \\ \begin{pmatrix} \frac{\partial H}{\partial S} \end{pmatrix}_{P} = T \qquad \begin{pmatrix} \frac{\partial H}{\partial P} \end{pmatrix}_{S} = V \\ \begin{pmatrix} \frac{\partial A}{\partial T} \end{pmatrix}_{V} = -S \qquad \begin{pmatrix} \frac{\partial A}{\partial V} \end{pmatrix}_{T} = -P \\ \begin{pmatrix} \frac{\partial G}{\partial T} \end{pmatrix}_{P} = -S \qquad \begin{pmatrix} \frac{\partial G}{\partial P} \end{pmatrix}_{T} = V$$

Basic Thermodynamic Relations 2

Maxwell Relations

| From (6) | $\left(\frac{\partial T}{\partial V}\right)_{S} = -\left(\frac{\partial P}{\partial S}\right)_{V}$ | (10) |
|----------|--|------|
| From (7) | $\left(\frac{\partial T}{\partial P}\right)_{S} = \left(\frac{\partial V}{\partial S}\right)_{P}$ | (11) |
| From (8) | $\left(\frac{\partial S}{\partial V}\right)_T = \left(\frac{\partial P}{\partial T}\right)_V$ | (12) |
| From (9) | $-\left(\frac{\partial S}{\partial P}\right)_T = \left(\frac{\partial V}{\partial T}\right)_P$ | (13) |

Thermodynamic Equation of State

$$dU = TdS - PdV \qquad \dots (6)$$
$$\left(\frac{\partial U}{\partial V}\right)_T = T\left(\frac{\partial S}{\partial V}\right)_T - P$$

Substituting (12)

$$\left(\frac{\partial U}{\partial V}\right)_T = T\left(\frac{\partial P}{\partial T}\right)_V - P$$

$$\left(\frac{\partial U}{\partial V}\right)_T = T\left(\frac{\alpha}{\kappa}\right) - P$$

How Free Energy Depends on T

| G = H - TS | (1) | definition |
|------------|-----|------------|
|------------|-----|------------|

 $\Delta G = \Delta H - T \Delta S \qquad \dots (2) \qquad \text{constant T}$

dG = VdP - SdT ...(3) fundamental eqn.

| | $\left(\frac{\partial G}{\partial T}\right)_{P} = -S = \frac{G - H}{T}$ | from (1) |
|--------------------------|---|--------------------------|
| | $\left(\frac{\partial G}{\partial T}\right)_{P} - \frac{G}{T} = -\frac{H}{T}$ | (4) |
| But | $\left(\frac{\partial (G/T)}{\partial T}\right)_{P} = \frac{1}{T} \left(\frac{\partial G}{\partial T}\right)_{P} - \frac{G}{T^{2}}$ | (5) |
| Gibbs-Helmho Equation | oltz $\left(\frac{\partial (G/T)}{\partial T}\right)_P = -\frac{H}{T^2}$ | substitute (4) in (5) |

alternative form:

$$\left(\frac{\partial (G/T)}{\partial (1/T)}\right)_P = H$$

By applying the Gibbs-Helmholtz equation to both reactants and products of a chemical reaction,

$$\left(\frac{\partial(\Delta G/T)}{\partial T}\right)_{P} = -\frac{\Delta H}{T^{2}}$$

How Free Energy Depends on P

$$\left(\frac{\partial G}{\partial P}\right)_T = V$$
$$\Delta G = G_2 - G_1 = \int_1^2 V dP$$

For solids and liquids V does not change much with P, so

$$G_2(T) \approx G_1(T) + V(P_2 - P_1)$$

For a perfect gas

$$V = \frac{nRT}{P}$$

$$\Delta G_T = nRT \ln(P_2 / P_1)$$

$$G = G^{\circ}(T) + nRT\ln(P/P^{\circ})$$

where G° is the standard free energy defined at $P^{\circ}=1$ bar

The chemical potential for a pure substance is the molar Gibbs energy:

$$\mu = \overline{G} = \frac{G}{n}$$

: For a perfect gas

$$\mu = \mu^{\circ} + RT \ln(P/P^{\circ})$$

The Chemical Potential

So far, our thermodynamic relations apply to closed systems. However, *G*, *U*, *H*, etc. are extensive properties.

Partial molar quantities are used to describe quantities which depend on composition.

$$\overline{J}_{i} = \left(\frac{\partial J}{\partial n_{i}}\right)_{T,P,n_{j \neq i}} \qquad J = U, H, S, A, G, V$$

$$J = \sum_{i} n_{i} \overline{J}_{i}$$
In general $dG = \left(\frac{\partial G}{\partial P}\right)_{T,n_{1},n_{2}...} dP + \left(\frac{\partial G}{\partial T}\right)_{P,n_{1},n_{2}...} dT$

$$+ \left(\frac{\partial G}{\partial n_{1}}\right)_{T,P,n_{2}...} dn_{1} + \left(\frac{\partial G}{\partial n_{2}}\right)_{T,P,n_{1},n_{3}...} dn_{2} + ...$$
Chemical Potential $\mu_{i} = \left(\frac{\partial G}{\partial n_{i}}\right)_{P,T,n_{j \neq i}}$
Fundamental Equations $dU = TdS - PdV + \sum_{i} \mu_{i} dn_{i}$
 $dH = TdS + VdP + \sum_{i} \mu_{i} dn_{i}$
etc.
$$\Rightarrow \quad \mu_{i} = \left(\frac{\partial U}{\partial n_{i}}\right)_{V,S} = \left(\frac{\partial H}{\partial n_{i}}\right)_{P,S} = \left(\frac{\partial A}{\partial n_{i}}\right)_{V,T}$$

CHEM 360 Spring 2004

Free Energy of Mixing 1

Consider the mixing of two *ideal* gases A and B at constant T and P.

Before:
$$G_1 = n_A \mu_A (\text{pure}) + n_B \mu_B (\text{pure})$$

where $\mu_A (\text{pure}) = \mu_A^\circ + RT \ln(P/P_0)$
After mixing: $G_A = n_A \left\{ \mu_A^\circ + RT \ln(p_A/P_0) \right\}$

$$= n_{\rm A} \left\{ \mu_{\rm A}^{\circ} + RT \ln(P/P_0) + RT \ln \chi_{\rm A} \right\}$$
$$= n_{\rm A} \mu_{\rm A} (\text{pure}) + n_{\rm A} RT \ln \chi_{\rm A}$$
$$G_2 = G_{\rm A} + G_{\rm B} < G_1$$
negative

Change:

$$\Delta G_{\text{mix}} = G_2 - G_1$$

= $n_A RT \ln \chi_A + n_B RT \ln \chi_B$
= $nRT (\chi_A \ln \chi_A + \chi_B \ln \chi_B)$
$$\Delta G_{\text{mix}} = nRT (\chi_A \ln \chi_A + \chi_B \ln \chi_B)$$

 $\Delta G_{\text{mix}} < 0$ Mixing is spontaneous.

Entropy of Mixing $\Delta S_{\text{mix}} = -(\partial \Delta G_{\text{mix}} / \partial T)_{P,n_i} = -nR \sum \chi_i \ln \chi_i$ Enthalpy of Mixing $\Delta H_{\text{mix}} = \Delta G_{\text{mix}} + T\Delta S_{\text{mix}} = 0$ Similarly, $\Delta U_{\text{mix}} = \Delta V_{\text{mix}} = 0$ for ideal gases and solutions

Free Energy of Mixing 2

For a binary system, ΔG_{mix} is at its minimum and ΔS_{mix} at its maximum when $\chi_A = \chi_B = 0.5$



Equilibrium and Mixing:

Consider an equilibrium between two ideal gases $A \rightleftharpoons B$

$$G = G_{A}(\text{pure}) + G_{B}(\text{pure}) + \Delta G_{\text{mix}}$$
$$= n(\chi_{A}\mu_{A} + \chi_{B}\mu_{B}) + nRT(\chi_{A}\ln\chi_{A} + \chi_{B}\ln\chi_{B})$$



The equilibrium composition is determined by the minimum in G_{total} .

Reaction Equilibrium 1

Consider $A \rightleftharpoons B$

Suppose an amount dx of A turns into B.

Then
$$(dG)_{T,P} = \mu_A dn_A + \mu_B dn_B = -\mu_A dx + \mu_B dx$$

$$\Rightarrow \qquad \left(\frac{\partial G}{\partial x}\right)_{T,P} = \mu_B - \mu_A$$

 μ_A and μ_B depend on composition, and therefore change during the reaction.

If $\mu_A > \mu_B$ the reaction proceeds from $A \rightarrow B$. If $\mu_A < \mu_B$ the reaction proceeds from $B \rightarrow A$.



Reaction Equilibrium 2

A reaction such as $2A + 3B \rightarrow C + 2D$ can be written as

$$0 = -2A - 3B + C + 2D$$

or completely generally as $0 = \sum_{i} v_i R_i = 0$

Then

$$\left(\frac{\partial G}{\partial \xi}\right)_{T,P} = \sum_{i} v_{i} \mu_{i}$$

and at equilibrium,

$$\sum_{i} v_{i} \mu_{i} = 0$$

$$\sum_{i} \left(\nu_{i} \mu_{i}^{\circ} + \nu_{i} RT \ln \frac{p_{i}^{eq}}{P^{\circ}} \right) = 0$$
$$\Delta \overline{G}_{rxn}^{\circ} + RT \ln \left[\prod_{i} \left(\frac{p_{i}}{P^{\circ}} \right)_{eq}^{\nu_{i}} \right] = 0$$

$$\Delta \overline{G}_{\rm rxn}^{\circ} = -RT\ln K$$

$$K = \prod_{i} \left(\frac{p_i}{P^\circ}\right)_{\rm eq}^{v_i}$$